

Magneto-optics of massless and massive electrons

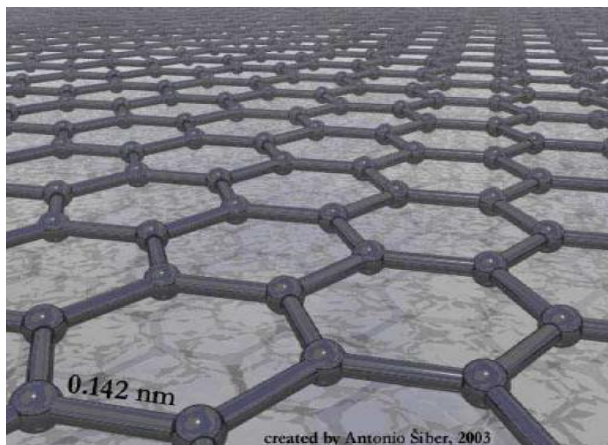
Milan Orlita

**Laboratoire National des Champs Magnétiques Intenses
CNRS, Grenoble, France**



Graphene

2D



2D crystal made of carbon atoms organized in hexagonal lattice

Theoretically known over sixty years...

P. R. Wallace, Phys. Rev. 71, 622 (1947)

Isolated/fabricated in 2004/2005

K. S. Novoselov et al., Science 306, 666 (2004)

K. S. Novoselov et al., Nature 438, 197 (2005)

3D

Diamond



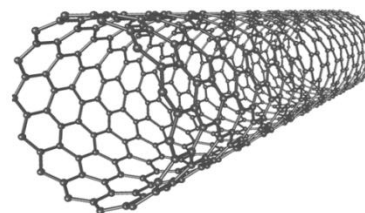
3D

Graphite



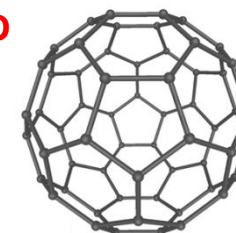
1D

Carbon nanotube



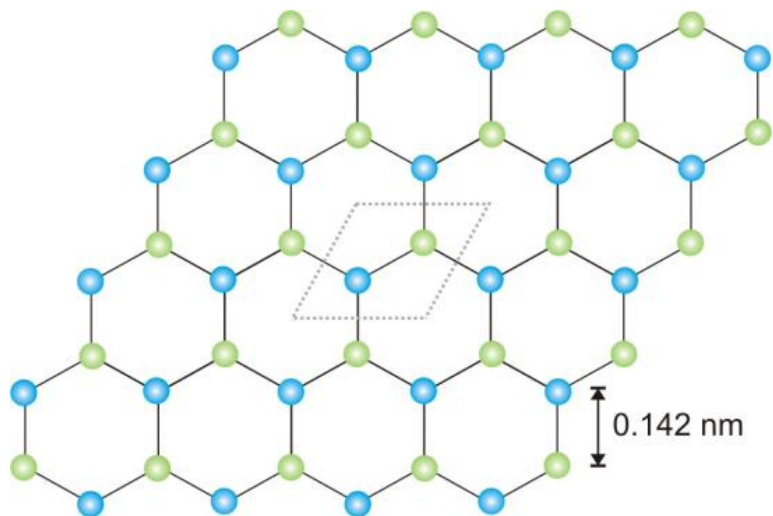
0D

Fulleren

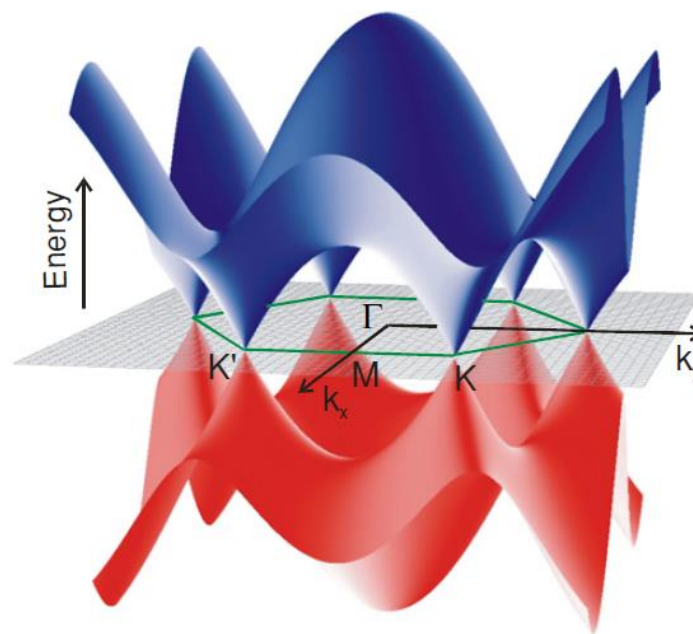


Electronic band structure of graphene

Crystal lattice:



Electronic bands:



Linearity of bands around K points:

$$E(\mathbf{k}) \approx \pm \hbar v |\mathbf{k}| = \pm v |\mathbf{p}|$$

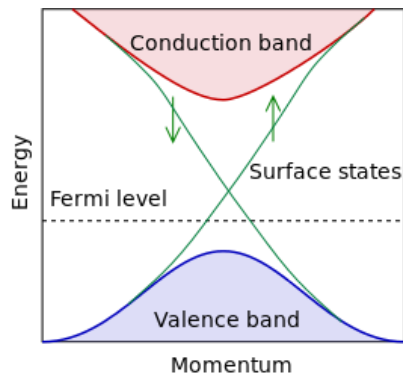
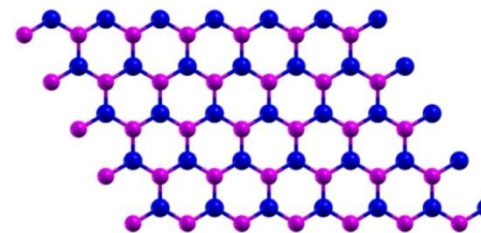
Electrons in graphene =
charged massless (relativistic) particles

$$E^2 = p^2 c^2 + \cancel{m_0^2 c^4}$$

Dirac-type materials: Dimension, symmetry, number of nodes...

Graphene, silicene, artificial graphene...

S. Cahangirov et al., Phys. Rev. Lett. 102, 236804 (2009)
 H. Liu et al., ACS Nano 8, 4033 (2014)
 C.-H. Park et al., PRL 101, 126804 (2008)
 M. Gibertini et al., Phys. Rev. B 79, 241406 (2009)



Topological and topological crystalline insulators

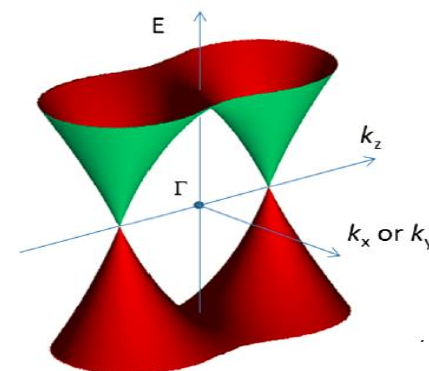
HgTe QWs, Bi_2Se_3 ,
 Bi_2Te_3 , $\text{Bi}_x\text{Sb}_{1-x}$, PbSnTe...

M. König et al., Science 318, 776 (2007)
 D. Hsieh et al., Nature 452, 970 (2008)
 H. Zhang et al., Nature Phys. 5, 438 (2009)
 L. Fu, Phys. Rev. Lett. 106, 106802 (2011)

3D Dirac and Weyl semimetals

Na_3Bi , Cd_3As_2 , TaAs, NbAs...

Z. K Liu et al., Science 343, 864 (2014)
 Z. K. Liu et al., Nature Mater. 13, 677 (2014)
 S. Jeon et al., Nature Mater. 13, 851 (2014)
 S. Borisenko et al., Phys. Rev. Lett. 113, 027603 (2014)
 M. Neupane et al., Nature Comm. 5, 3786 (2014)
 B. Q. Lv et al., Nature Phys. 11, 724 (2015)
 L. X. Yang et al., Nature Phys. 11, 728 (2015)
 S.-Y. Xu et al., Nature Phys. 11, 748 (2015)



Outline:

Magneto-optics of massive Dirac electrons in Bi_2Se_3

M. Orlita et al., Phys. Rev. Lett. 114, 186401 (2015)

C. Faugeras, B. A. Piot, G. Martinez,
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A. Akrap et al., arXiv:1604.00038 (2016)

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A. Akrap, Wednesday 1/6/16, 9:50

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Relativistic quantum electrodynamics in condensed-matter physics...

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Dirac Hamiltonian (4x4):

$$H_D = m_D c^2 \vec{\beta} + c \vec{\alpha} \cdot \vec{p}$$

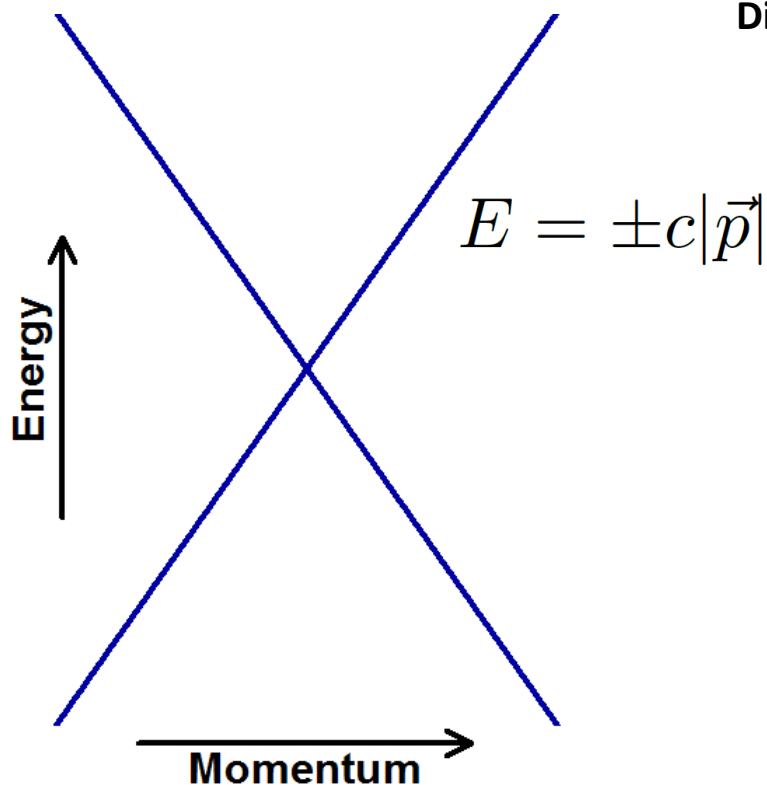
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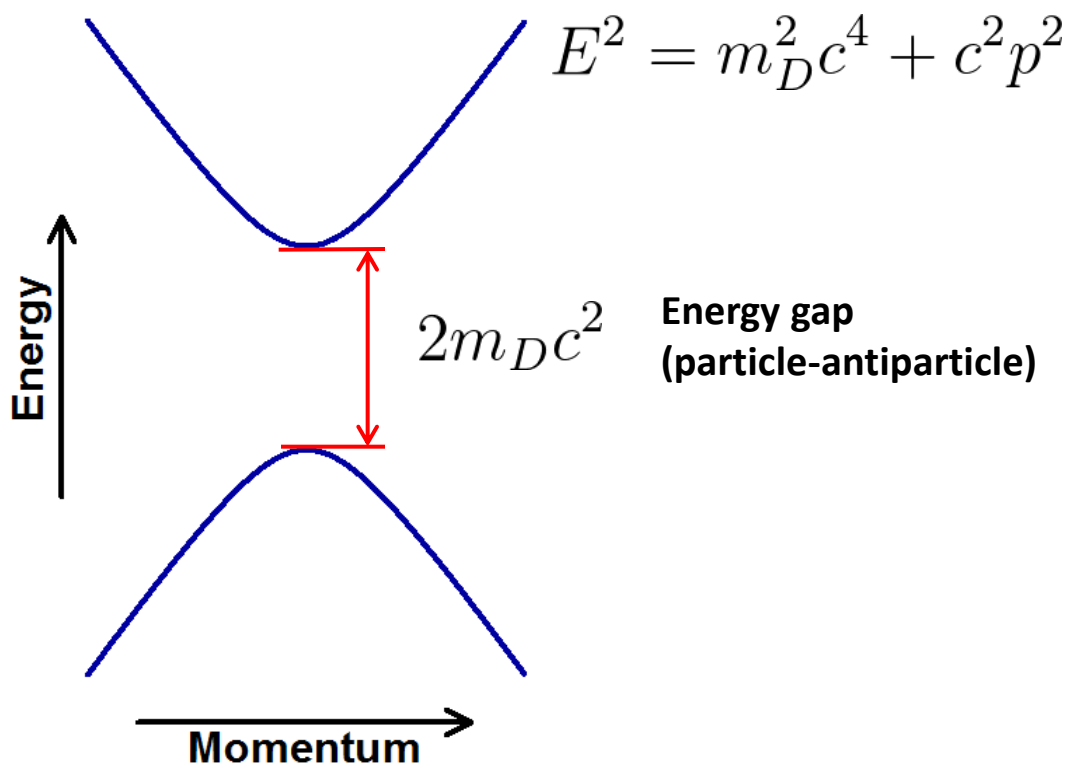
Dirac mass (= rest mass)

$$m_D = 0$$



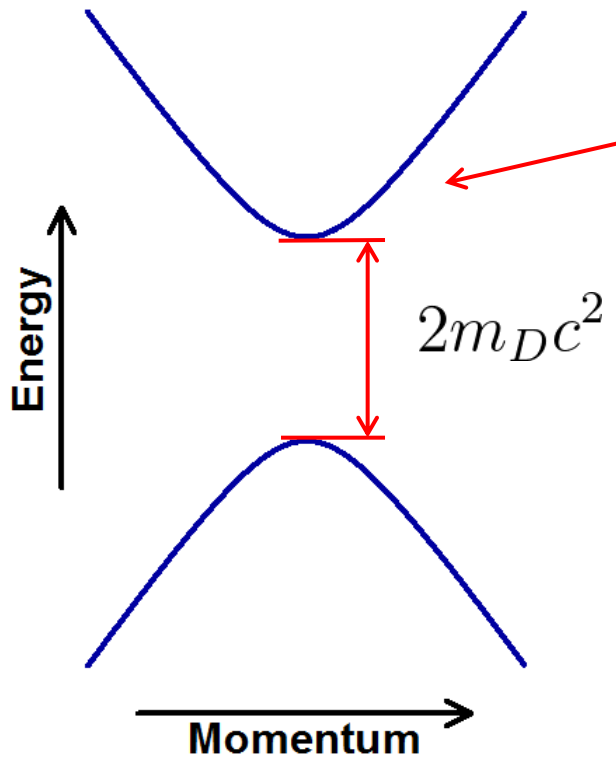
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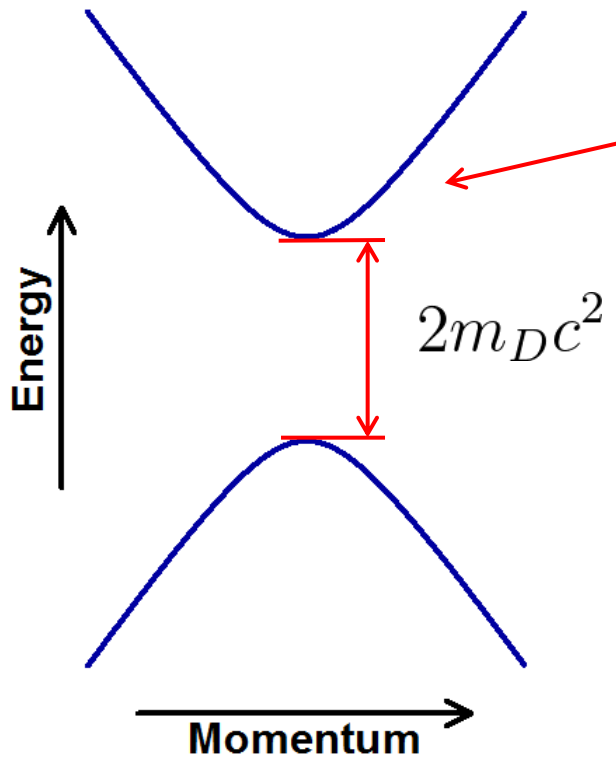
Particle mass m_D
(in the classical limit)

Energy gap, particles masses,
and g factors determined by
two parameters...

$$m_D, c$$

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Dirac Hamiltonian (4x4):
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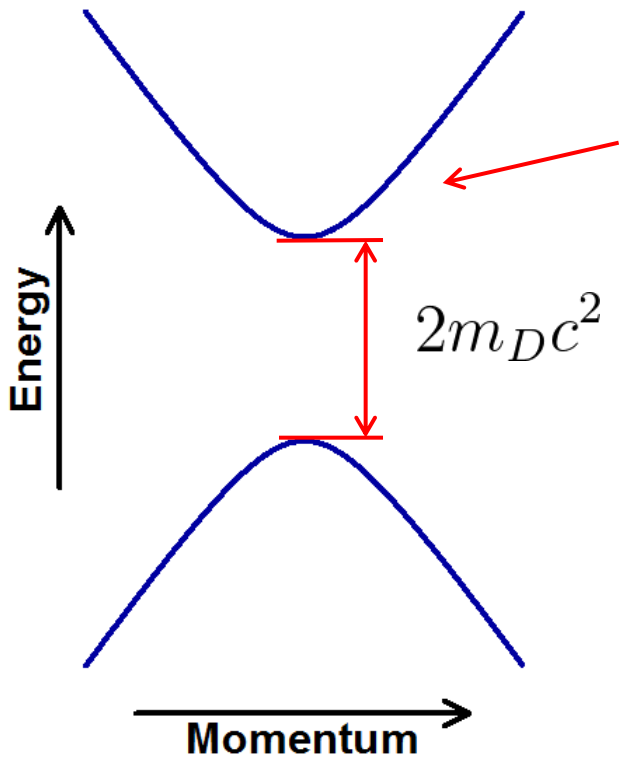
Particle mass m_D
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g factor (particle & antiparticle)

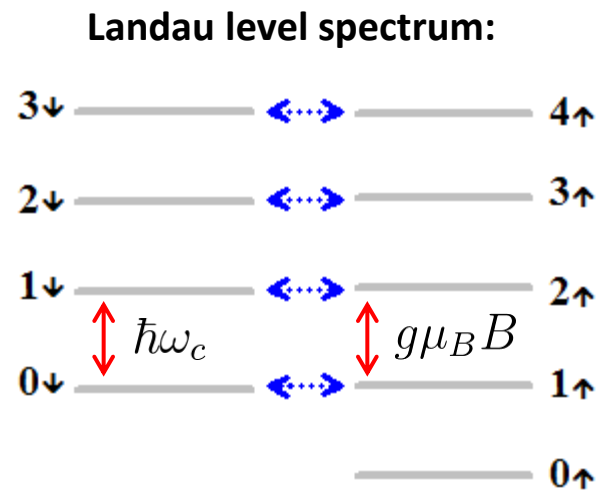
$$g = 2 \frac{m_0}{m_D}$$

Relativistic quantum electrodynamics in condensed-matter physics...

Dirac Hamiltonian (4x4):
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Particle mass m_D
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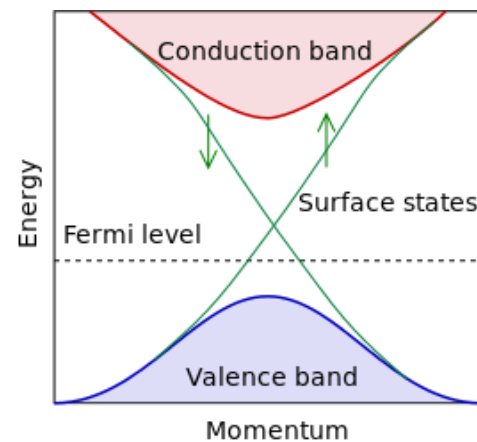


cyclotron energy = spin-splitting

Electronic bands in 3D topological insulators (Bi_2Se_3 family)

Topological insulators = narrow gap semiconductors
with conducting Dirac-type surface states

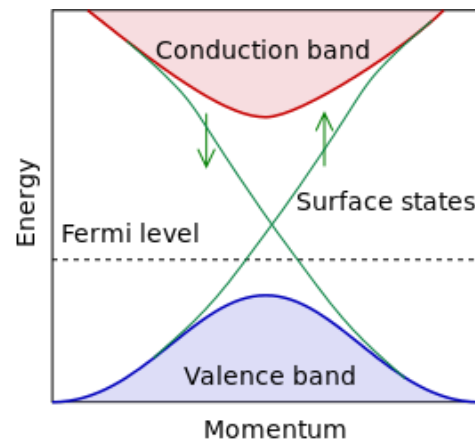
Bi_2Se_3 , Bi_2Te_3 , Sb_2Te_3 , $\text{Bi}_x\text{Sb}_{1-x}$



Electronic bands in 3D topological insulators (Bi₂Se₃ family)

Topological insulators = narrow gap semiconductors with conducting Dirac-type surface states

Bi₂Se₃, Bi₂Te₃, Sb₂Te₃, Bi_xSb_{1-x}

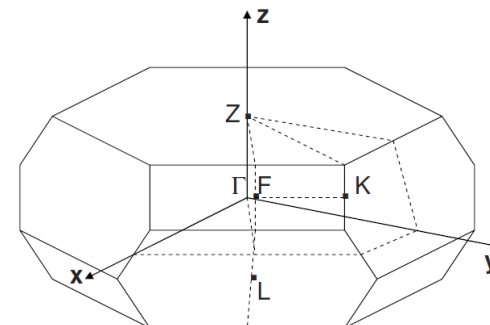


Effective 3D Dirac Hamiltonian at the Γ point:

$$H_D = Cp^2 \vec{I} + (m_D v_D^2 + Mp^2) \vec{\beta} + v_D \vec{\alpha} \cdot \vec{p}$$

\uparrow
e-h asymmetry

\uparrow
band inversion

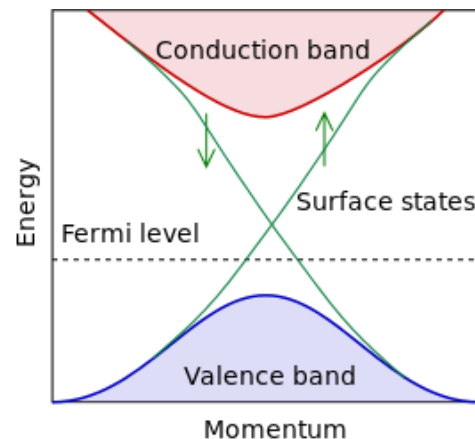


H. Zhang et al., Nature Phys. 5, 438 (2009)
C.-X. Liu et al., Phys. Rev. B 82, 045122 (2010)

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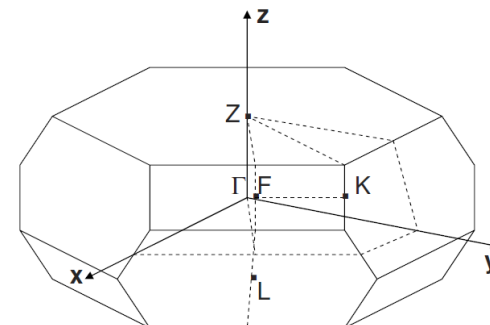


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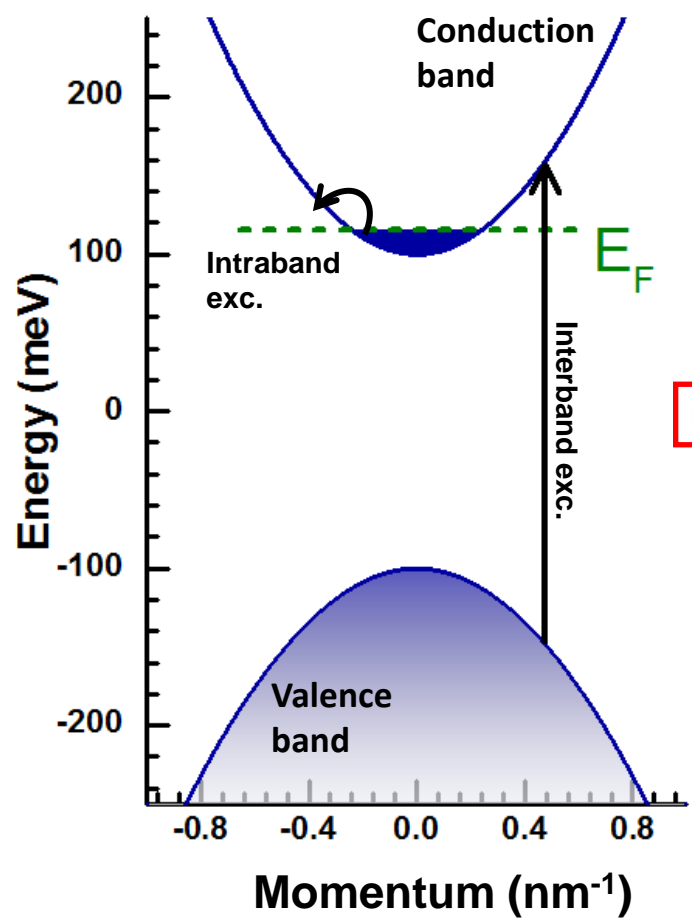
$$H_D = Cp^2 \vec{I} + (m_D v_D^2 + Mp^2) \vec{\beta} + v_D \vec{\alpha} \cdot \vec{p}$$

Electronic bands (electron and hole branches):

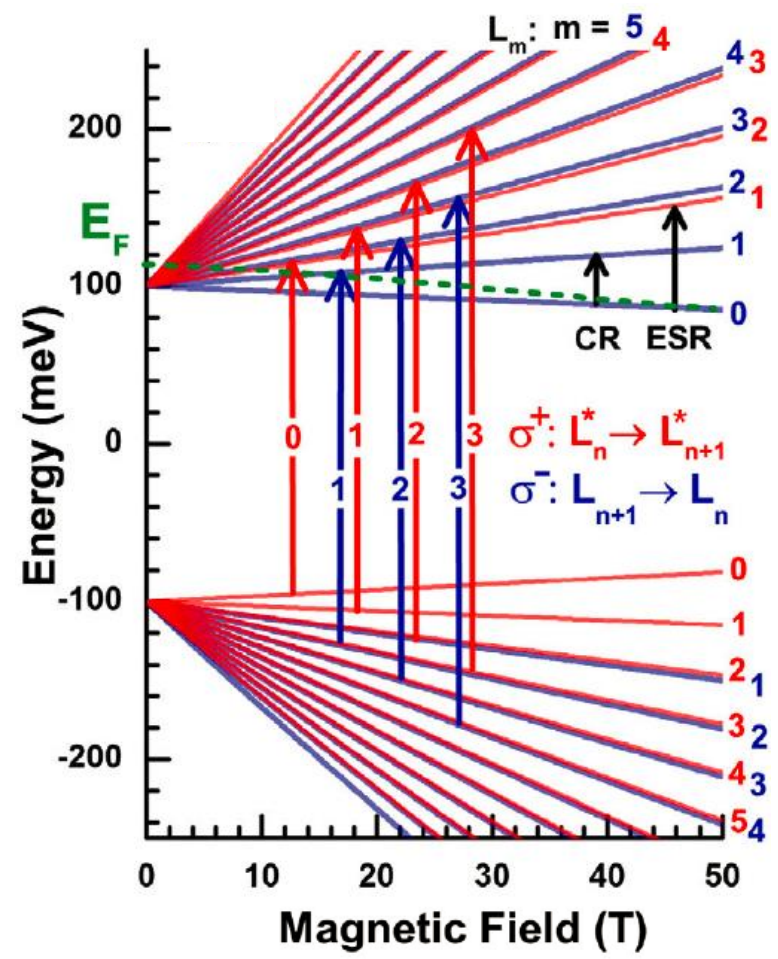
$$E = Cp^2 \pm \sqrt{(m_D v_D^2 + Mp^2)^2 + v_D^2 p^2}$$



Landau level spectroscopy

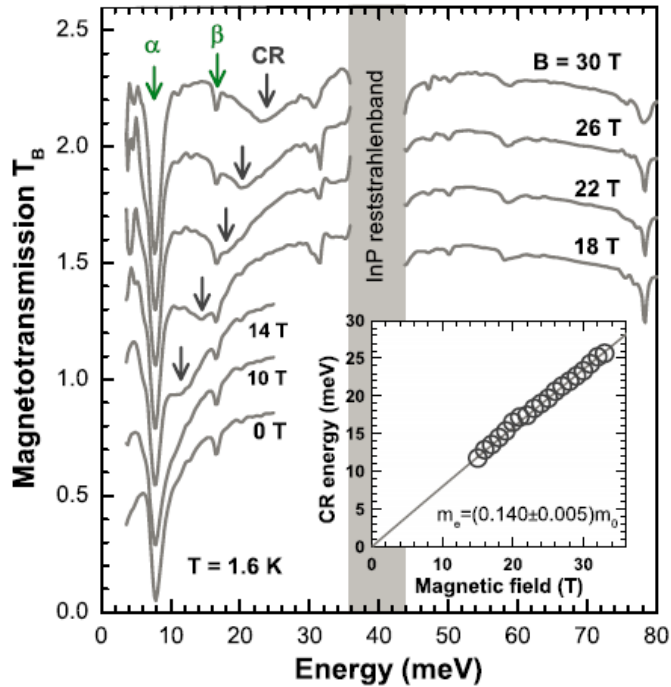


Magnetic field

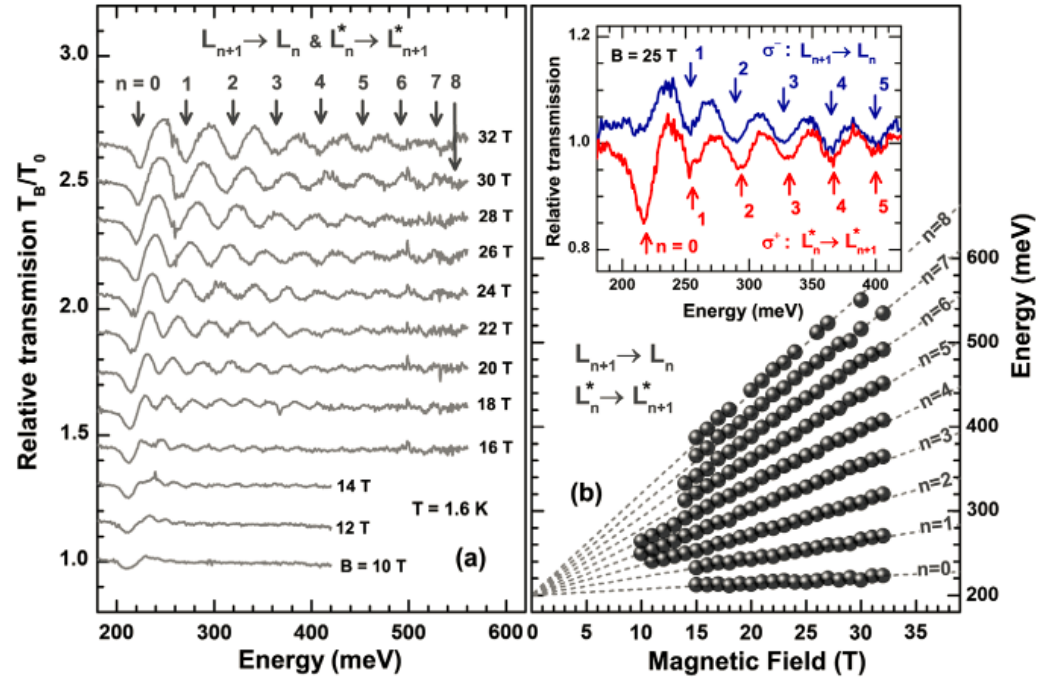


Thin layer of Bi_2Se_3 : Infrared magneto-transmission

Cyclotron resonance:



Interband inter-Landau level excitations:



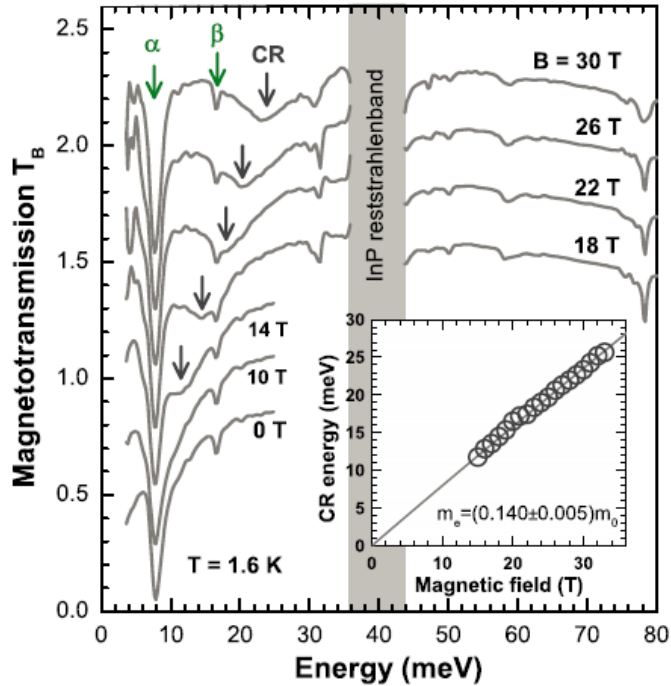
Bi_2Se_3 on InP(B) substrate
thickness 270 nm, $n \sim 5 \times 10^{17} \text{ cm}^{-2}$

Magneto-transmission experiment
on thin MBE-grown layers of Bi_2Se_3

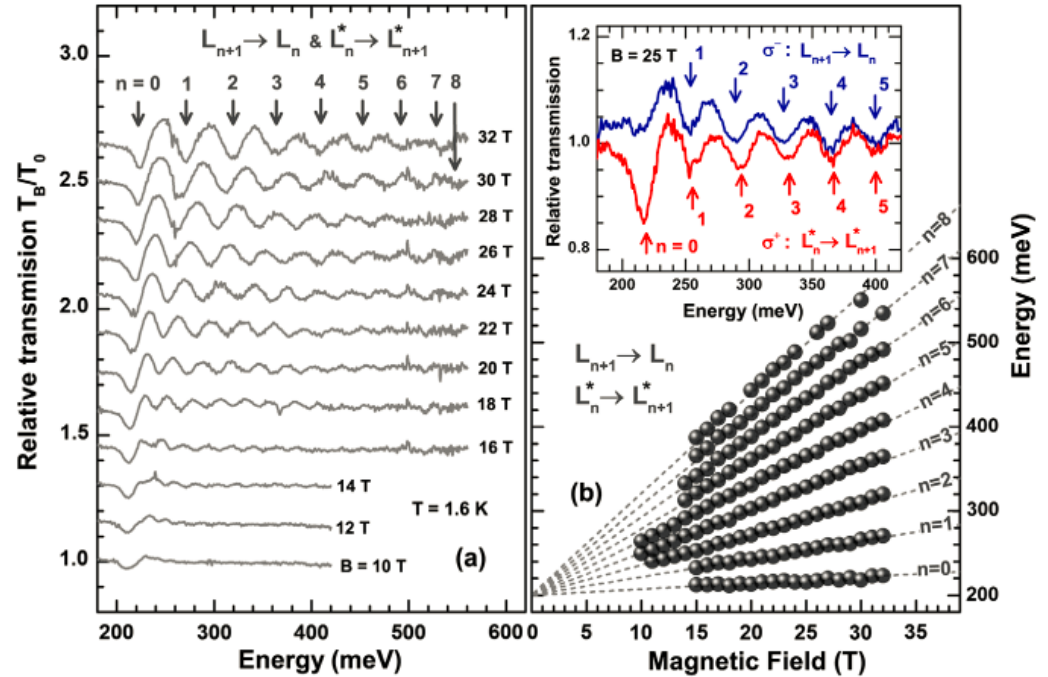
Group of L. W. Molenkamp

Thin layer of Bi₂Se₃: Infrared magneto-transmission

Cyclotron resonance:



Interband inter-Landau level excitations:



Bi₂Se₃ on InP(B) substrate
thickness 270 nm, $n \sim 5 \times 10^{17} \text{ cm}^{-2}$

Conduction and valence band parabolic

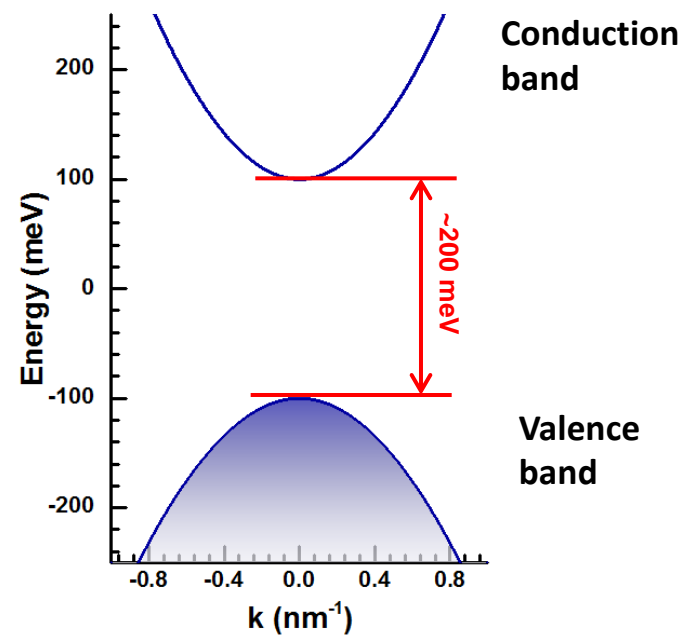
Low electron hole asymmetry $m_e \approx m_h$

Band gap $\sim 200 \text{ meV}$

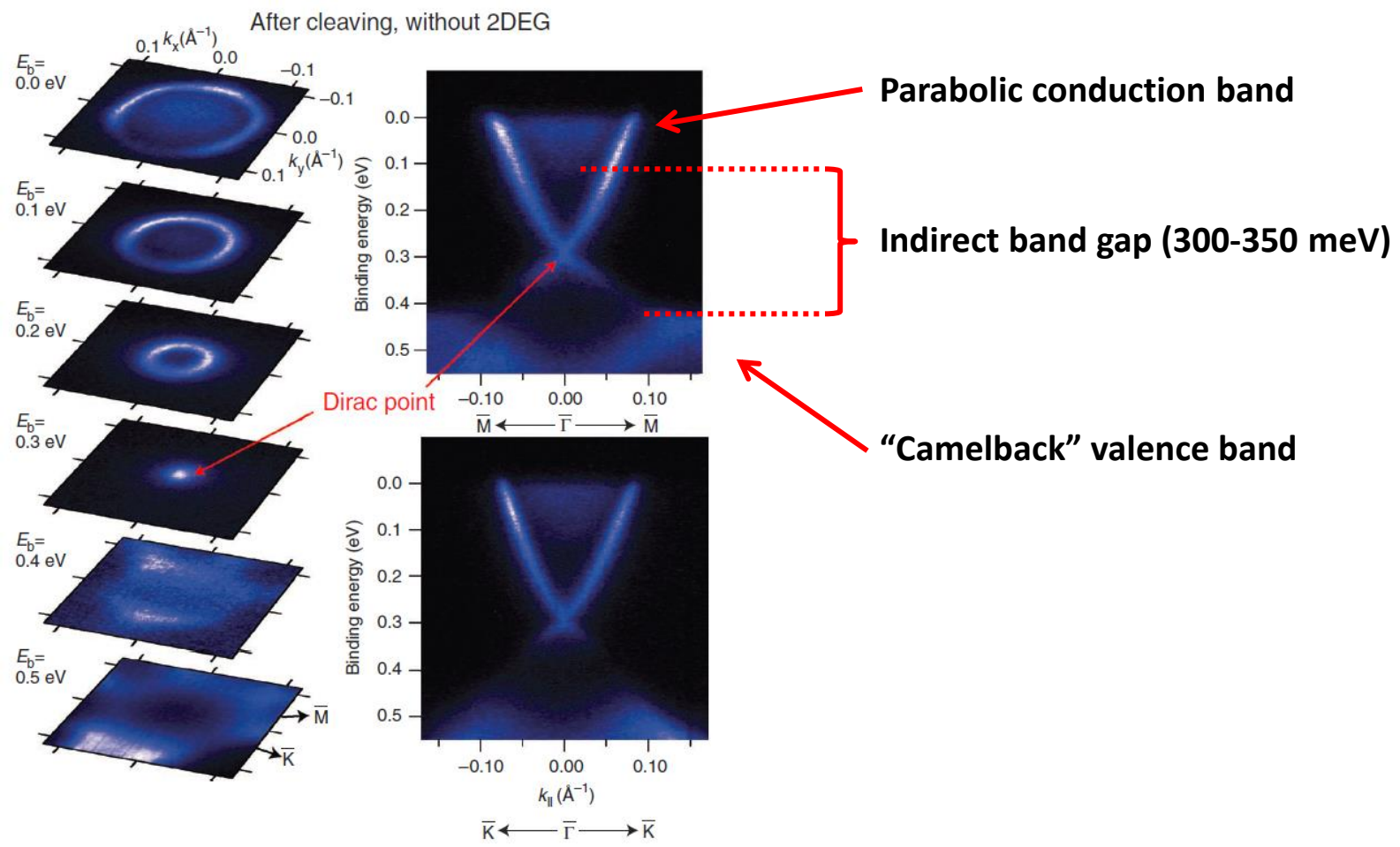
Electronic bands in bulk Bi_2Se_3 (Γ point)

Experiment (magneto-optics):

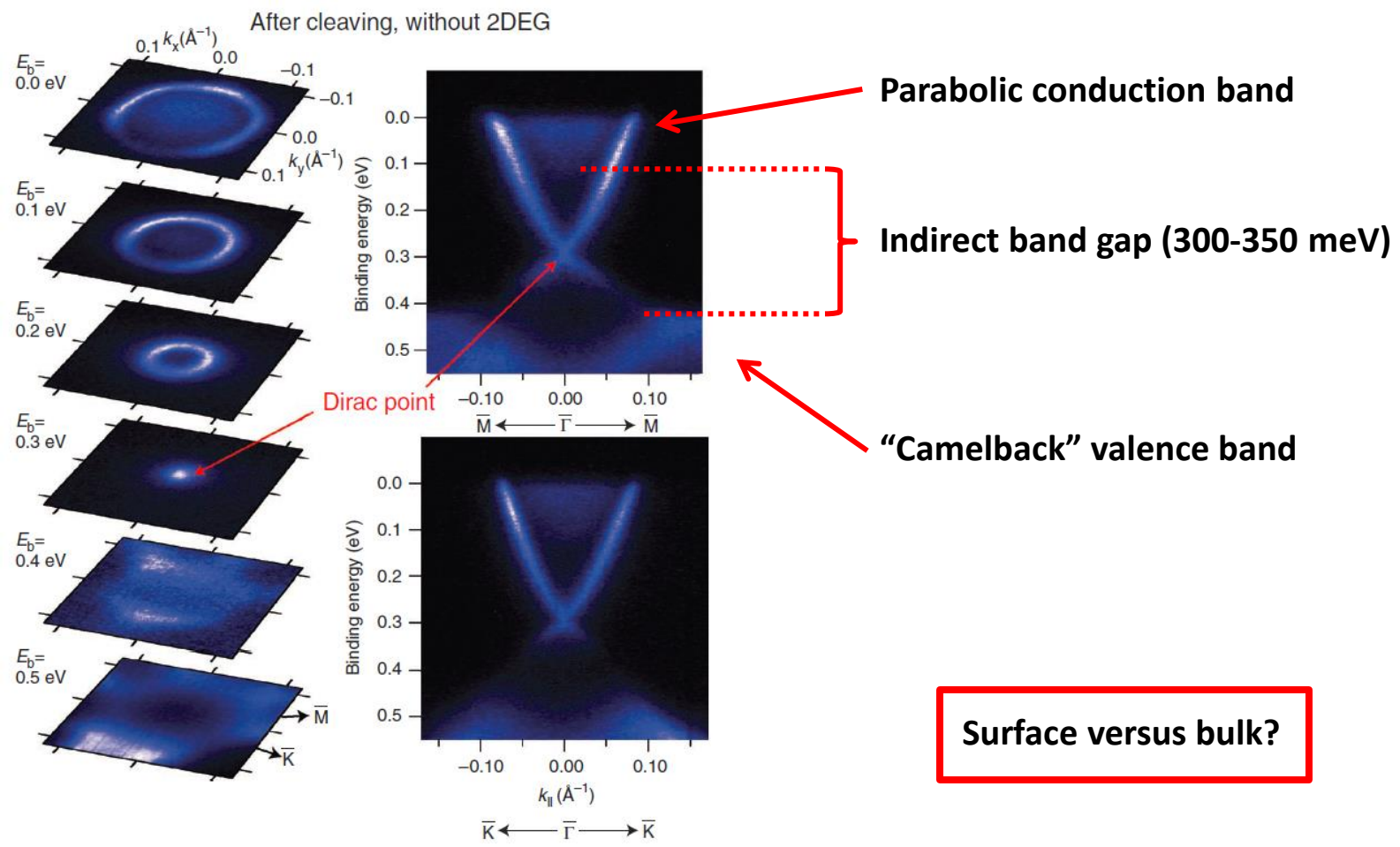
Conduction and valence band parabolic
High electron-hole symmetry



Digression: Band structure of Bi_2Se_3 in ARPES



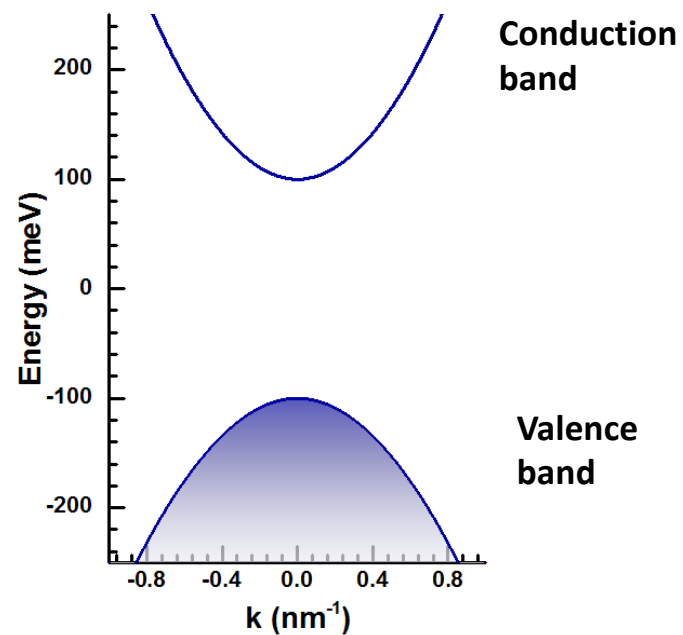
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Conduction and valence band parabolic
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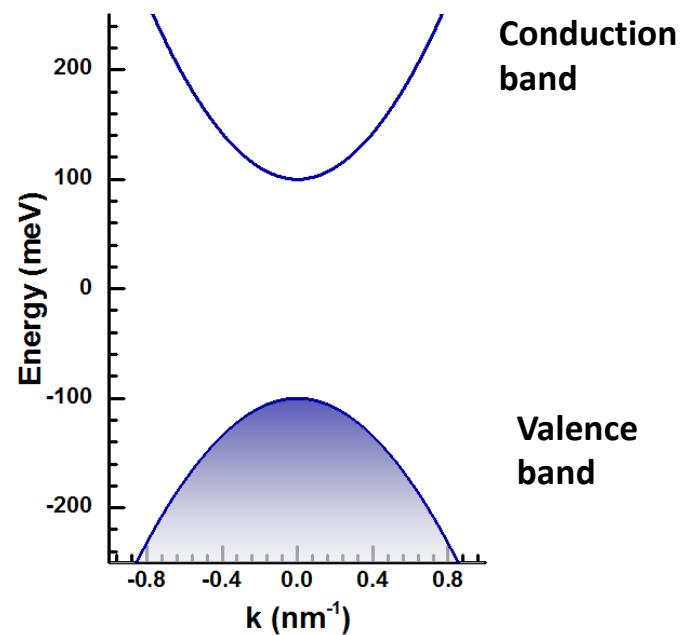
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Theory:

$$E = Ck^2 \pm \sqrt{(m_D v_D^2 + Mk^2)^2 + v_D^2 \hbar^2 k^2}$$

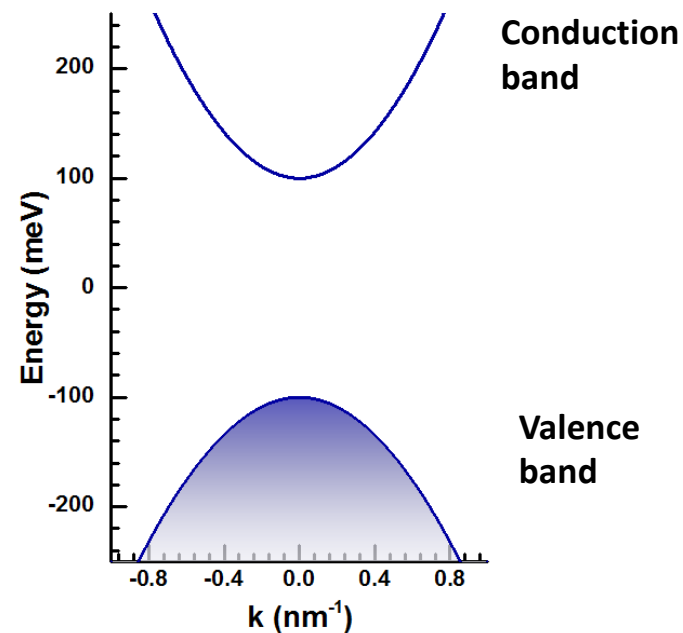


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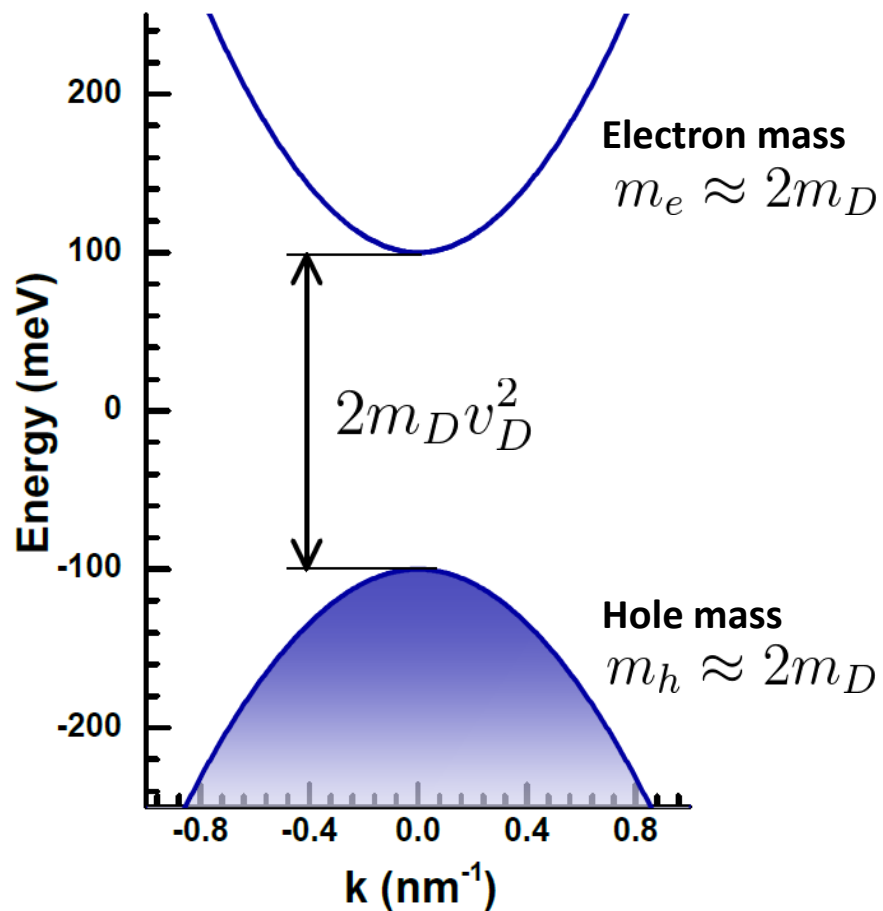
Theory:

$$E = \cancel{Ck^2} \pm \sqrt{(m_D v_D^2 + Mk^2)^2 + v_D^2 \hbar^2 k^2}$$

$$-\frac{\hbar^2}{4m_D} = M \quad \uparrow$$

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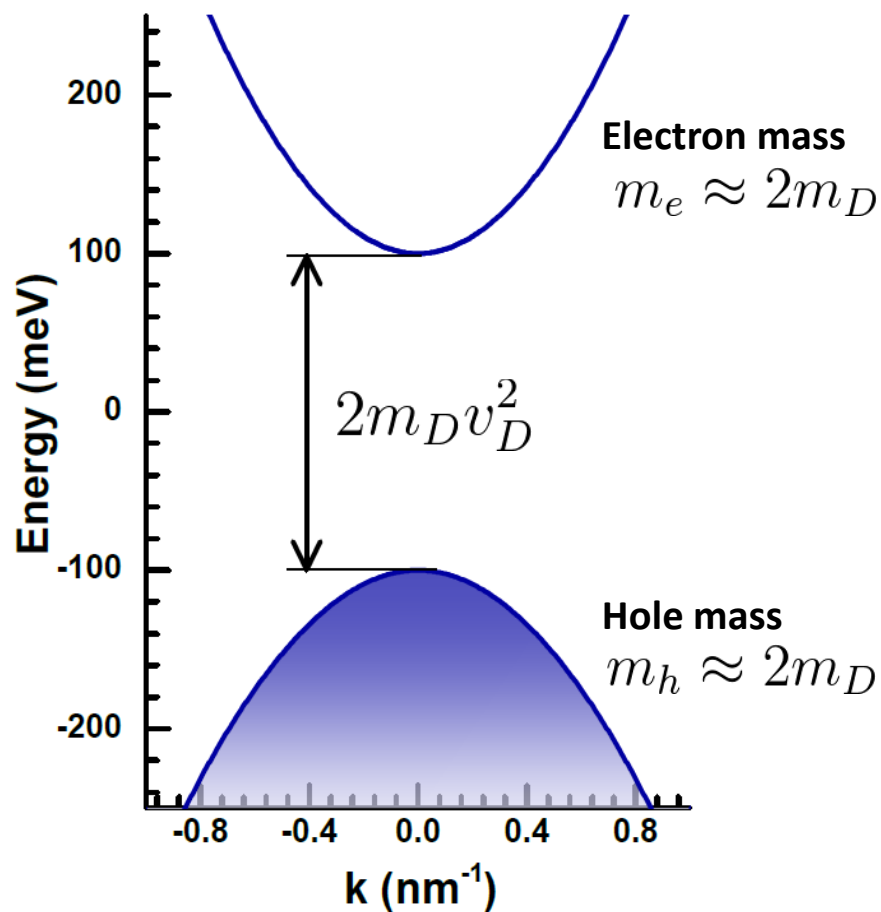
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Energy gap, particles masses and g factors
in Bi_2Se_3 determined by only two parameters...

$$m_D, v_D$$

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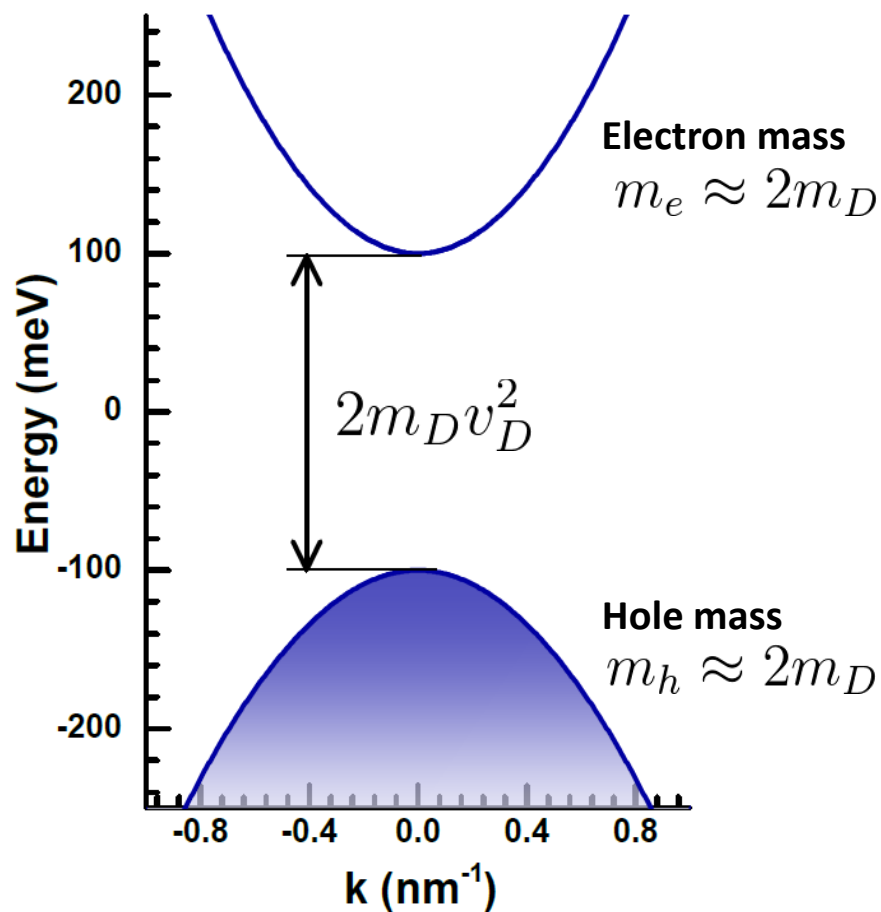
$$m_D, v_D$$

g factor:

$$g_e \approx g_h \approx 2 \frac{m_0}{m_D} \approx 25$$

(for $m_D \approx 0.08m_0$)

Electronic bands in bulk Bi_2Se_3 (Γ point)



Energy gap, particles masses and g factors
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$$m_D, v_D$$

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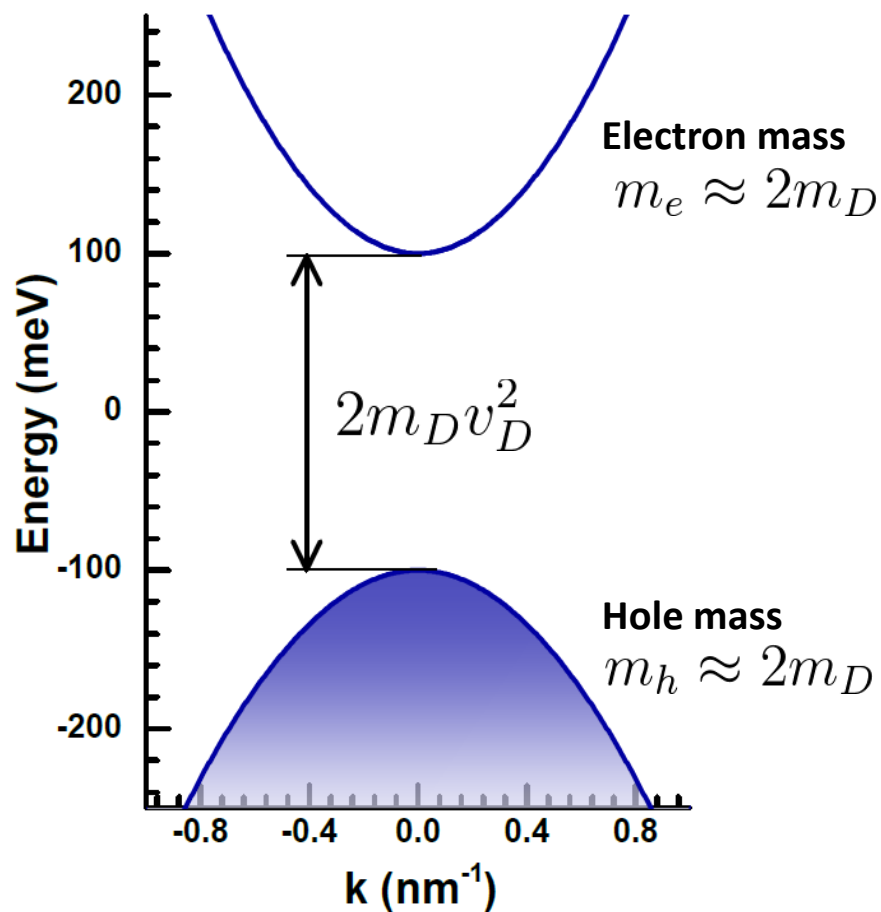
$$g_e \approx g_h \approx 2 \frac{m_0}{m_D} \approx 25$$

(for $m_D \approx 0.08m_0$)

Electron g factor from EPR:

$$g^{\text{ESR}} = 27.5$$

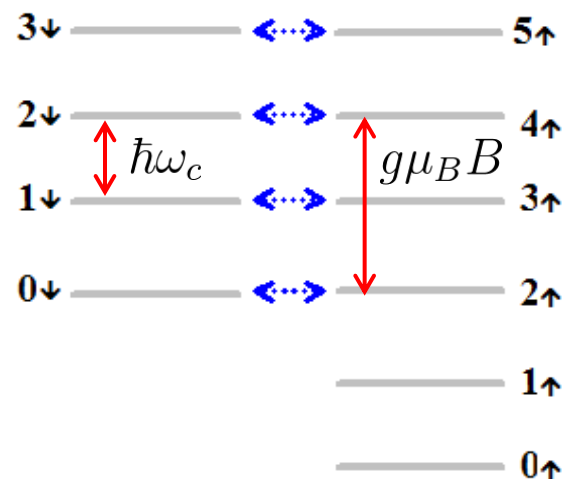
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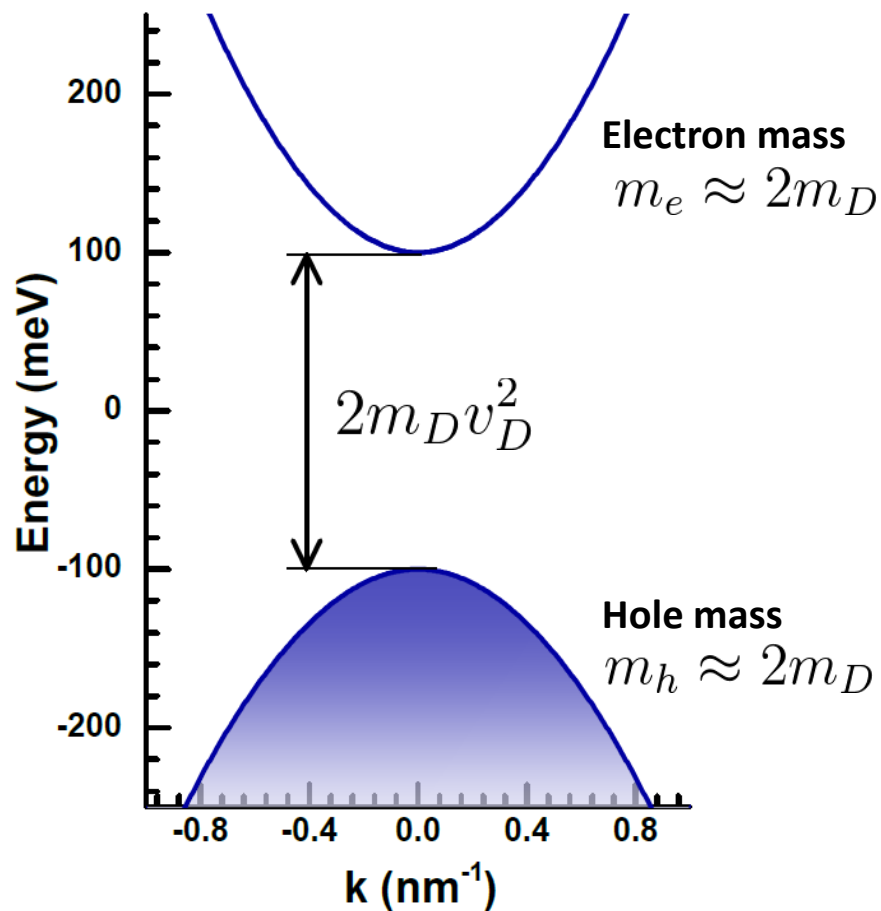
$$m_D, v_D$$

Landau level spectrum:



2× cyclotron energy = spin-splitting

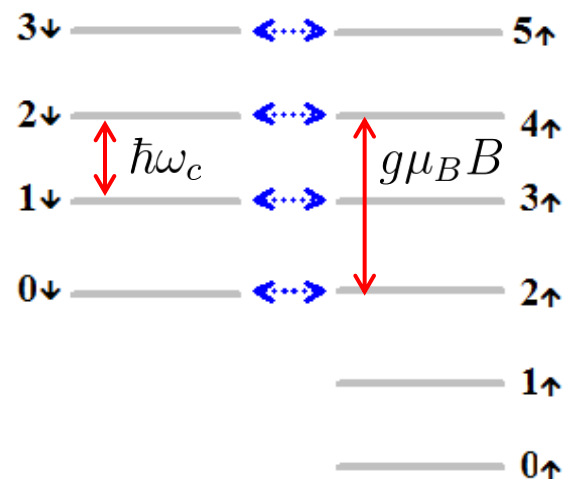
Electronic bands in bulk Bi_2Se_3 (Γ point)



Energy gap, particles masses and g factors
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$$m_D, v_D$$

Landau level spectrum:



Well-known empirical fact from magneto-transport...

see e.g. H. Köhler and E. Wöchner, phys. stat. sol. (b) 67, 665 (1975)

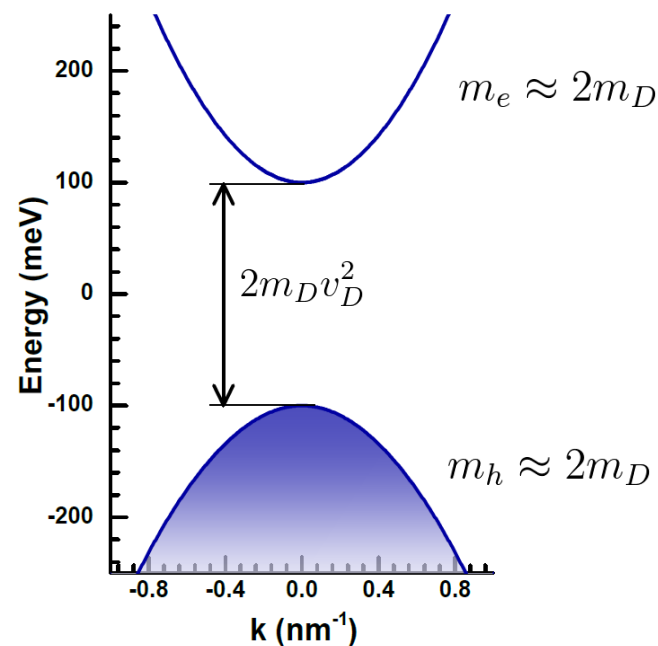
Electronic bands in bulk Bi_2Se_3 : Conclusions

Electrons and holes in bulk Bi_2Se_3 closely resemble massive Dirac particles in quantum electrodynamics

Energy gap, particles masses and g factors in Bi_2Se_3 determined by only two parameters...

$$m_D, v_D$$

In very good agreement with magneto-transport and EPR, but not ARPES.



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L. Ohnoutek et al., Sci. Rep. 6, 19087 (2016)

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Dirac semimetal Cd_3As_2 = stable 3D analogue of graphene

nature
materials

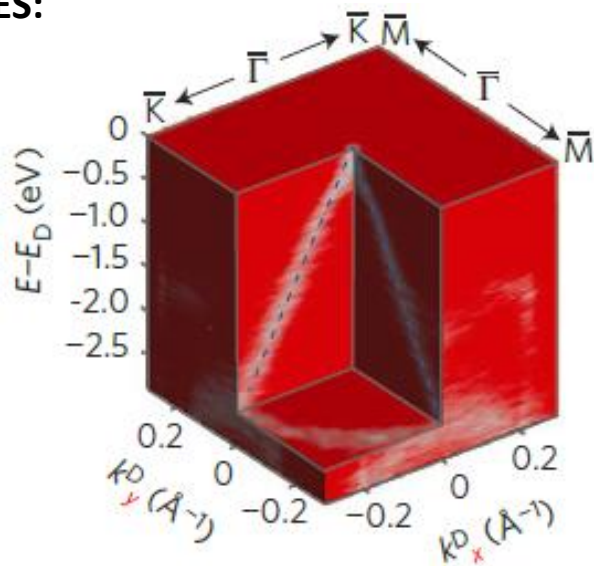
LETTERS

PUBLISHED ONLINE: 25 MAY 2014 | DOI: 10.1038/NMAT3990

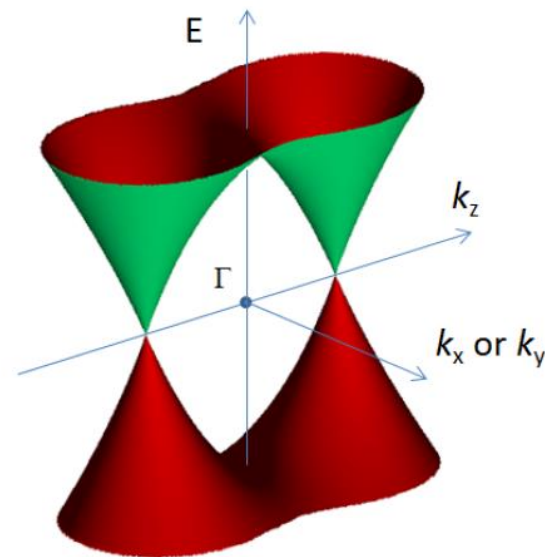
A stable three-dimensional topological Dirac semimetal Cd_3As_2

Z. K. Liu^{1†}, J. Jiang^{2,3†}, B. Zhou^{2,4†}, Z. J. Wang^{5†}, Y. Zhang^{1,4}, H. M. Weng⁵, D. Prabhakaran², S-K. Mo⁴, H. Peng², P. Dudin⁶, T. Kim⁶, M. Hoesch⁶, Z. Fang⁵, X. Dai⁵, Z. X. Shen¹, D. L. Feng³, Z. Hussain⁴ and Y. L. Chen^{1,2,4,6*}

ARPES:



3D Dirac semimetal Cd_3As_2



Z. K. Liu et al., Nature Mater. 13, 677 (2014)

S. Borisenko et al., Phys. Rev. Lett. 113, 027603 (2014)

M. Neupane et al., Nature Comm. 5, 3786 (2014)

Dirac semimetal Cd_3As_2 = stable 3D analogue of graphene

nature
materials

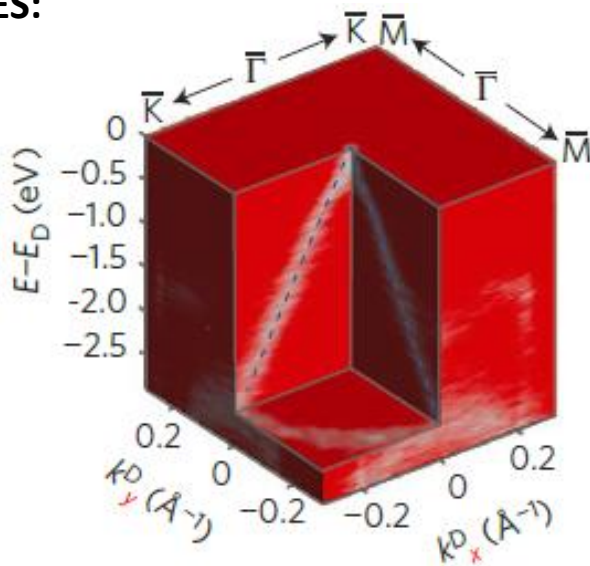
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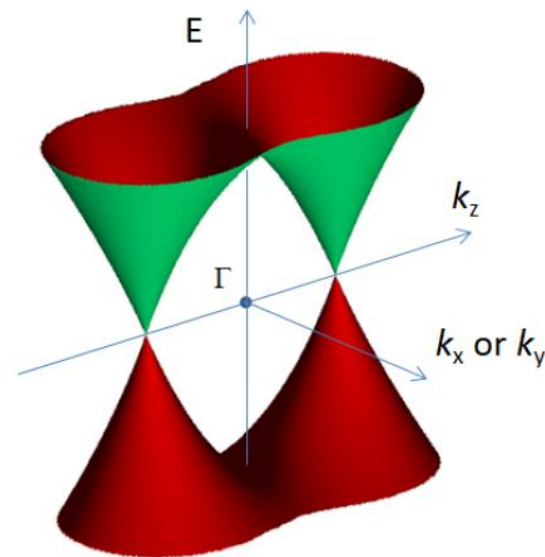
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ARPES:



3D Dirac semimetal Cd_3As_2

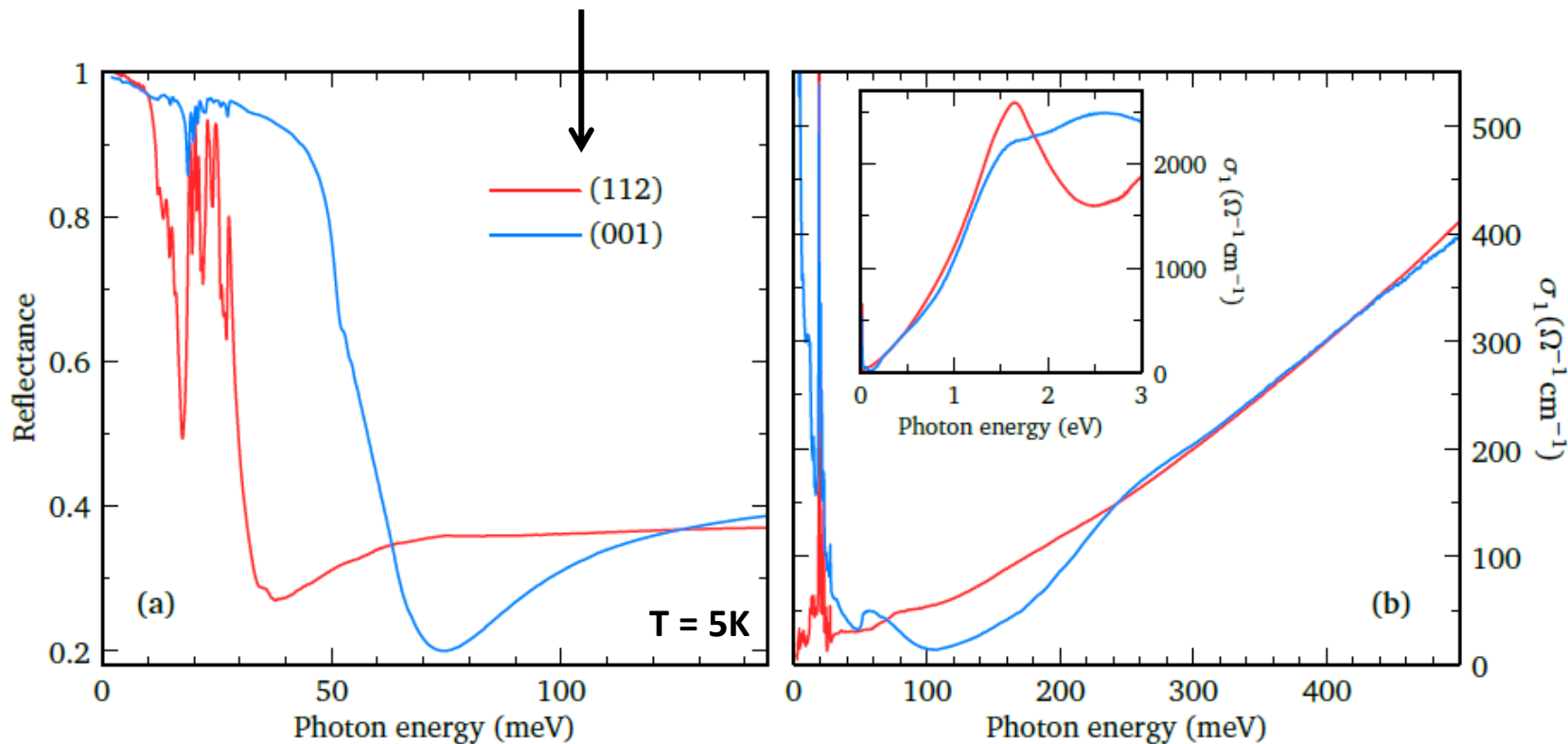


...but really not in line with other studies,
e.g. recent STS/STM experiments

S. Jeon et al., Nature Mater. 13, 851 (2014)

Cd_3As_2 – Infrared reflectance at $B=0$

Two different crystallographic orientations of **tetragonal** Cd_3As_2



...in collaboration with C. C. Homes (Brookhaven)

Cd₃As₂ – Infrared reflectance at B=0

Absorption of light in solids
(e.g., Fermi's golden rule):

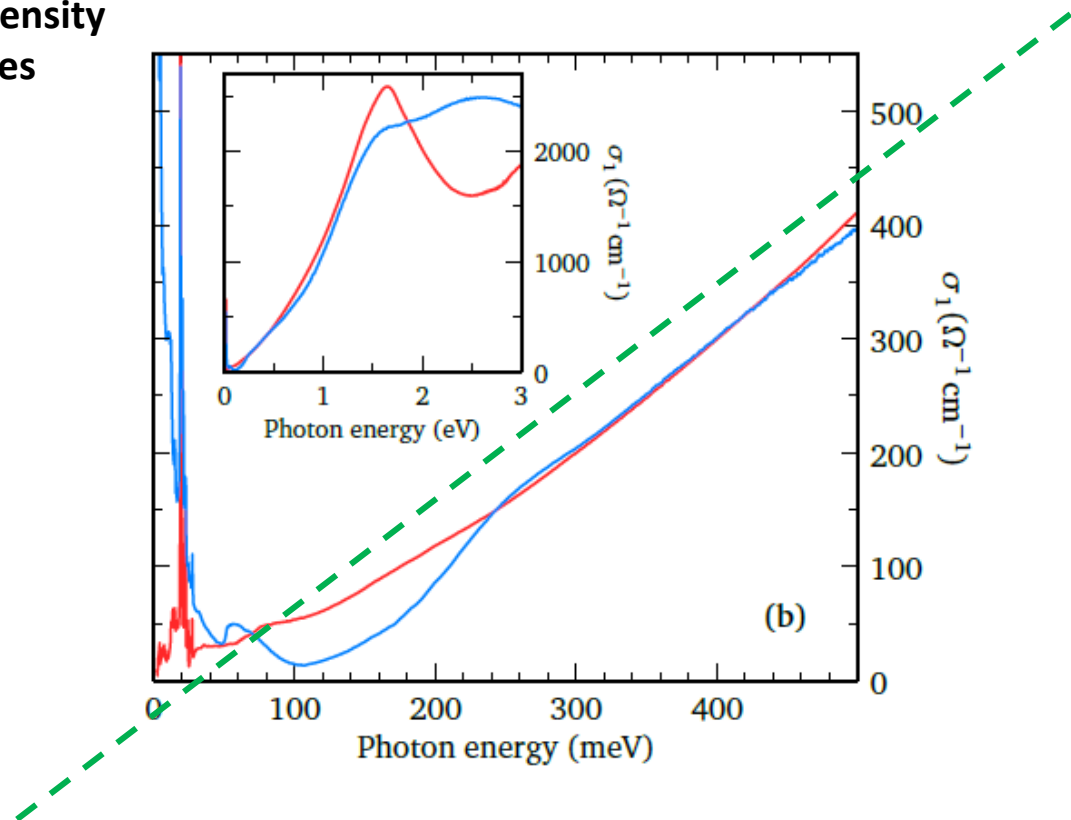
$$\sigma_1(\omega) \propto \frac{D(\omega)}{\omega}$$

joint density
of states

For conical bands in 3D:

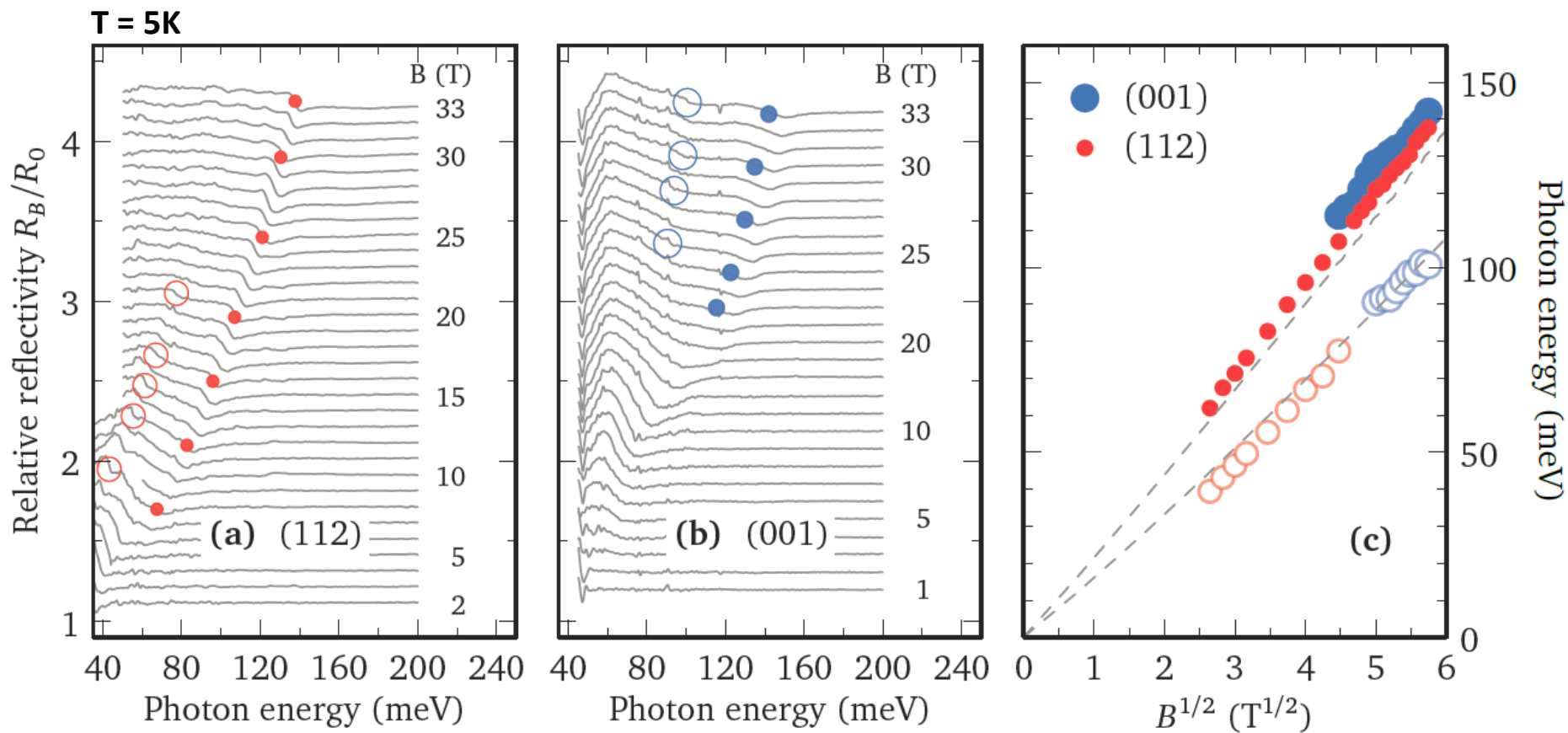
$$D(\omega) \propto \omega^2$$

Absorption coefficient (optical
conductivity) linear in photon
frequency!



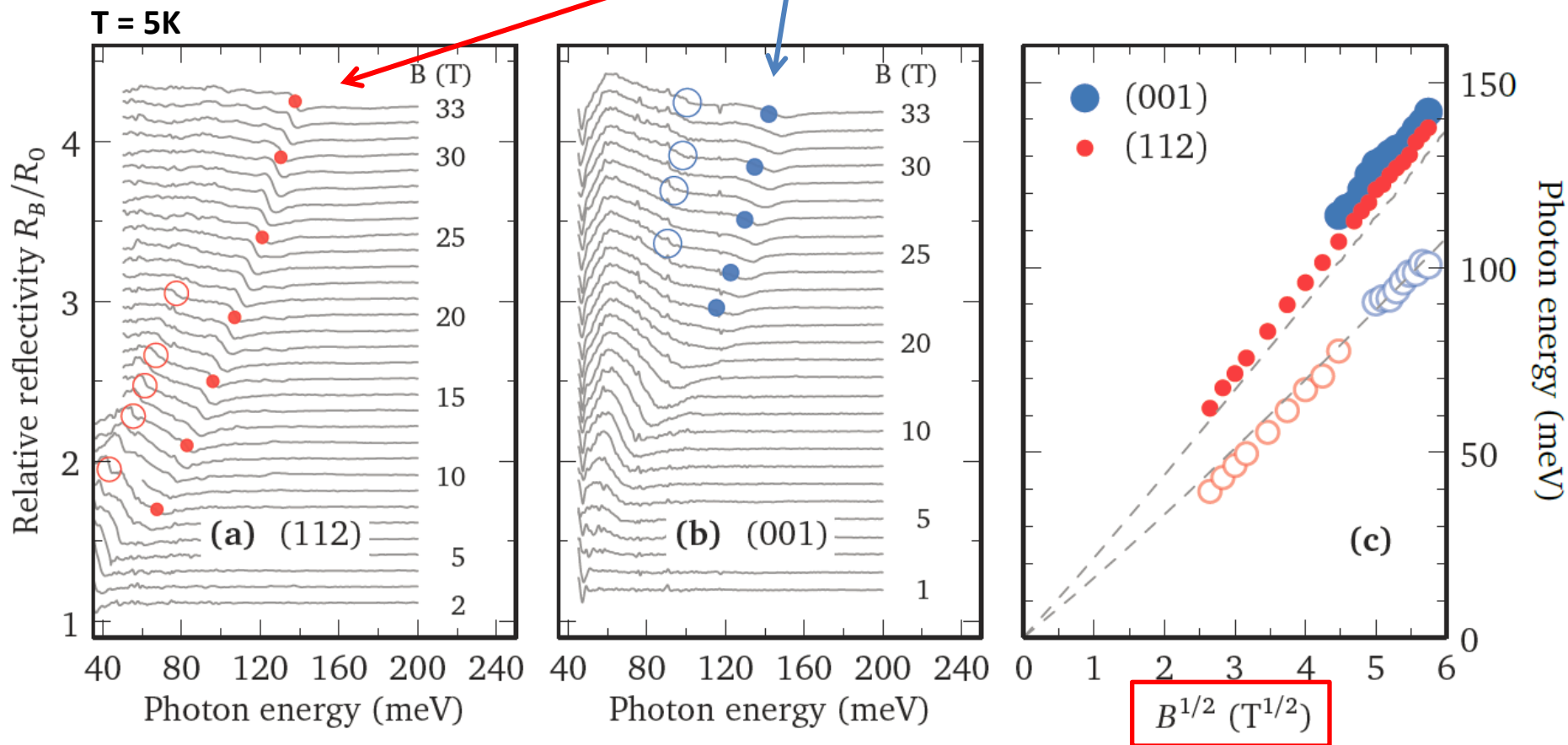
see also D. Neubauer et al. , Phys. Rev. B 93, 121202 (2016)

Cd₃As₂ – High-field magneto-reflectivity



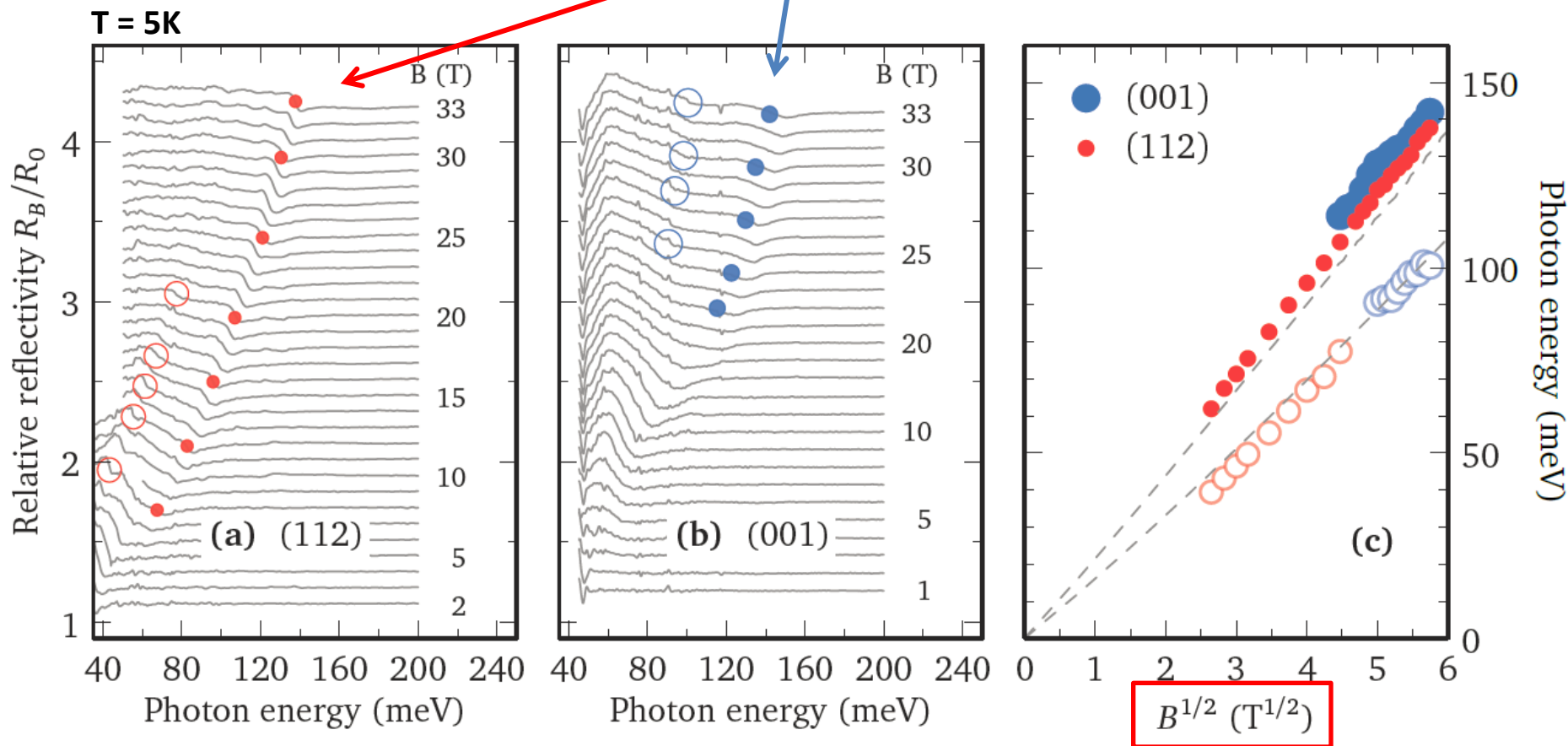
Cd₃As₂ – High-field magneto-reflectivity

Cyclotron resonance (CR) in the quantum limit



Cd₃As₂ – High-field magneto-reflectivity

Cyclotron resonance (CR) in the quantum limit

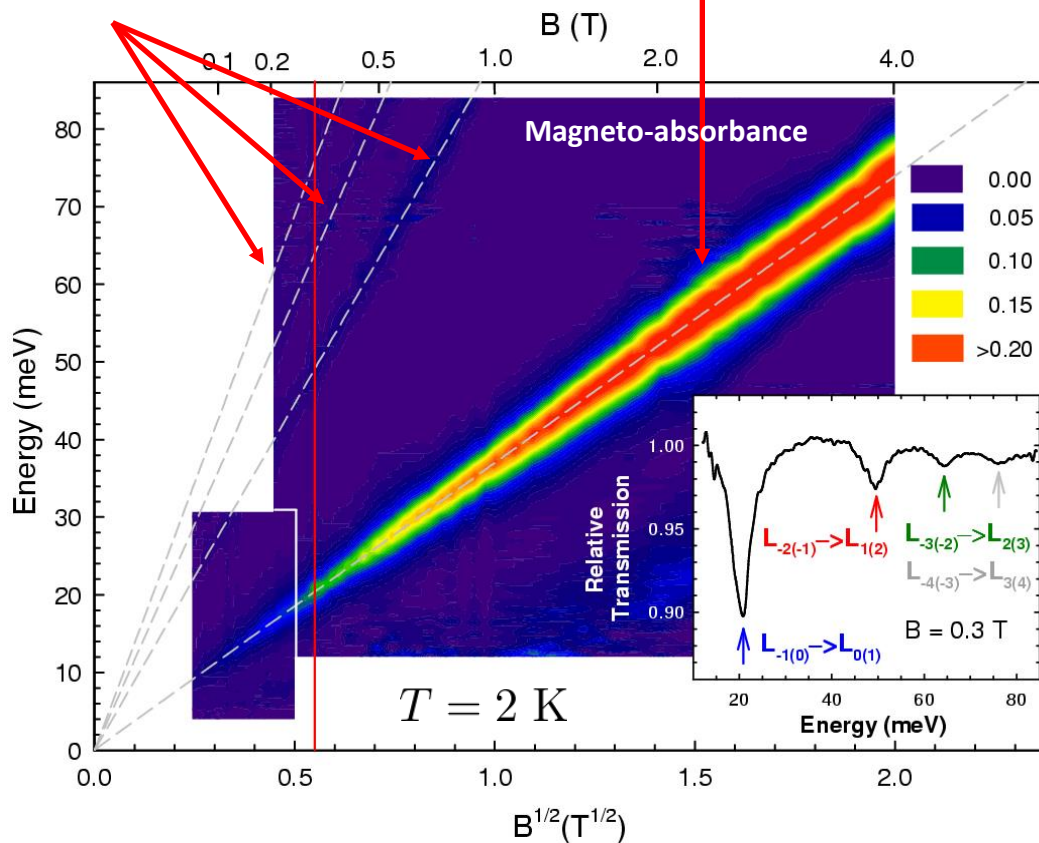


Magneto-optical response linear in \sqrt{B} = typical signature of massless particles

Magneto-transmission of (multilayer epitaxial) graphene

Interband inter-Landau level transitions

Cyclotron resonance



Energy spectrum:

$$E_n = \pm v \sqrt{2e\hbar B |n|}$$

Velocity:

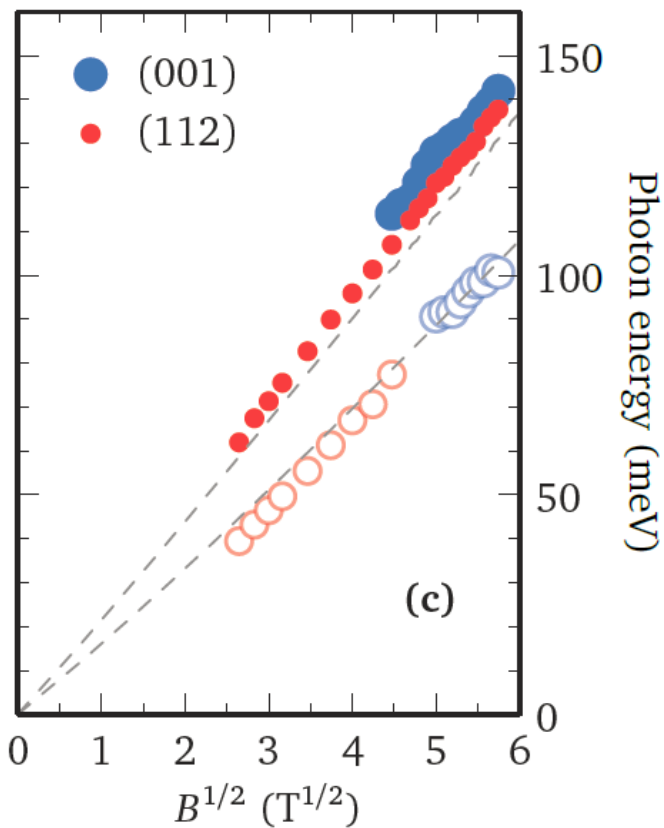
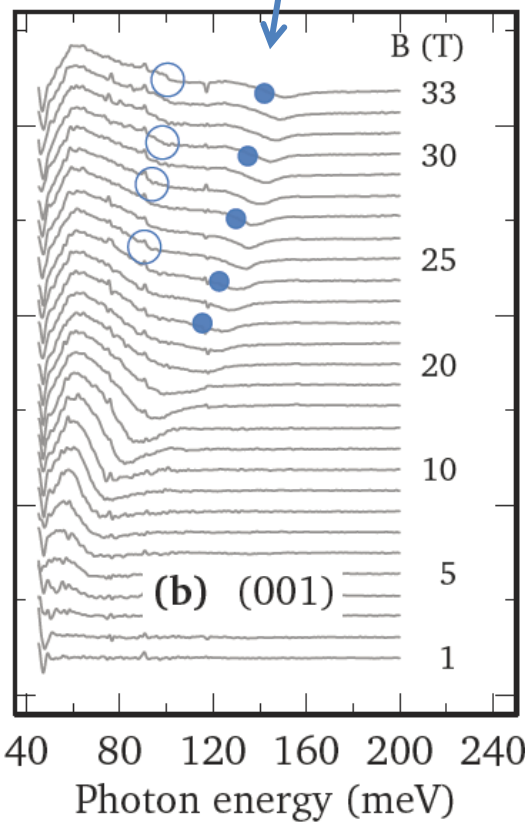
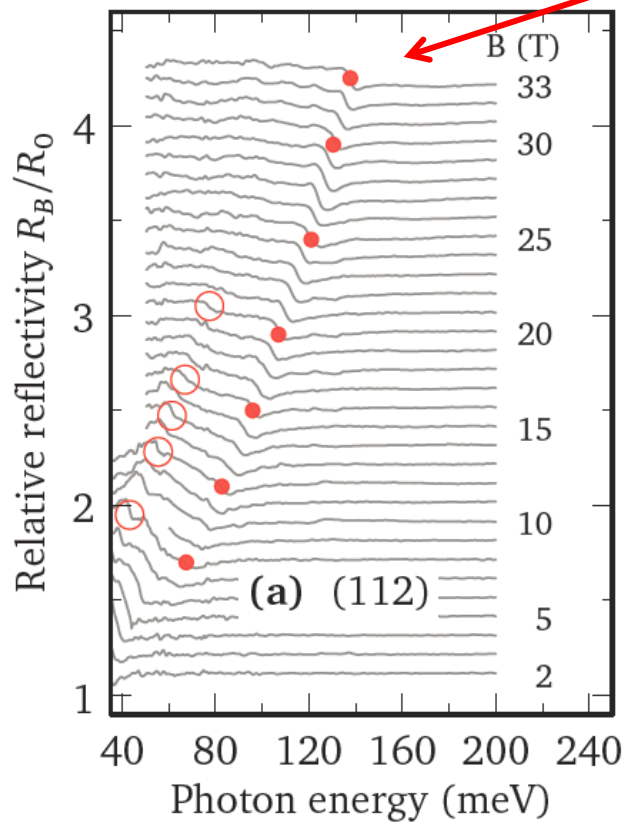
$$v = 1.02 \times 10^6 \text{ m/s}$$

Selection rules:

$$|n| \rightarrow |n| \pm 1$$

Cd₃As₂ – High-field magneto-reflectivity

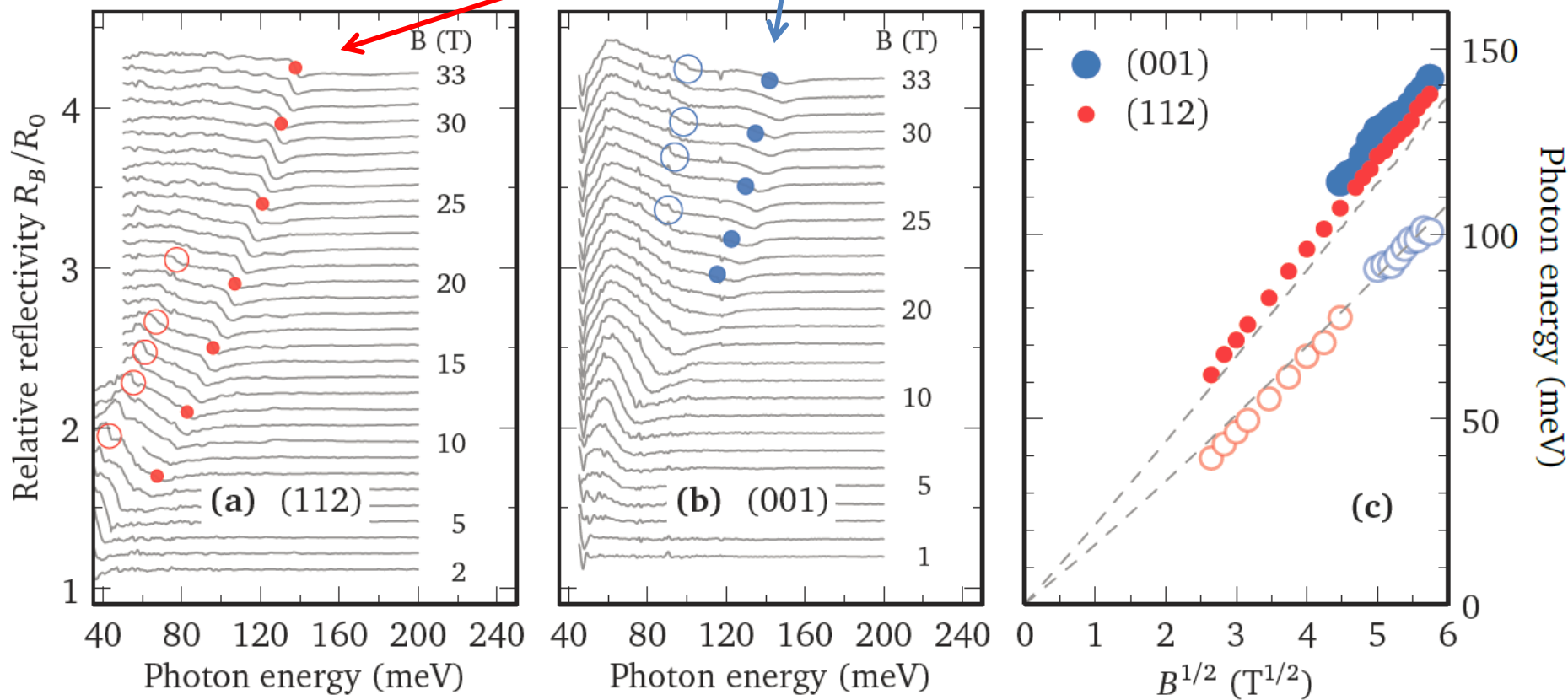
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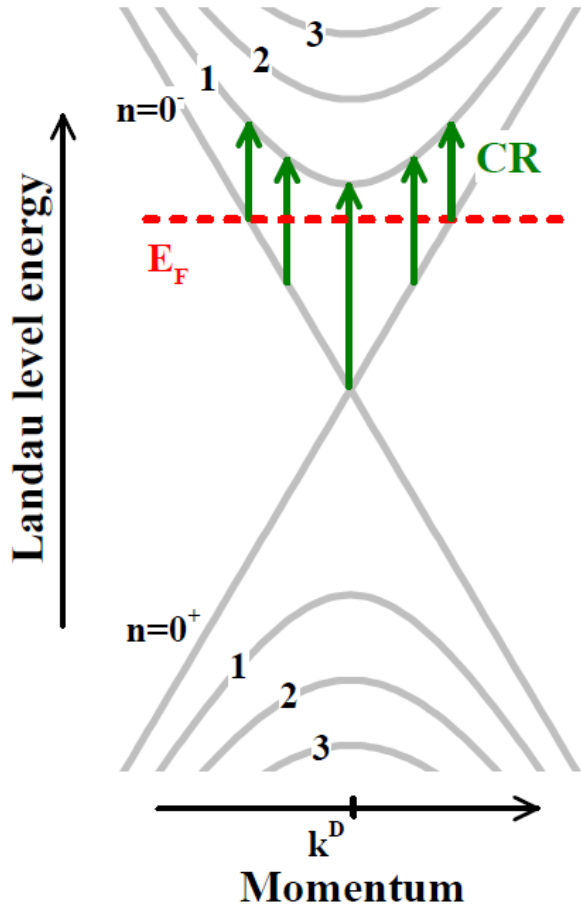


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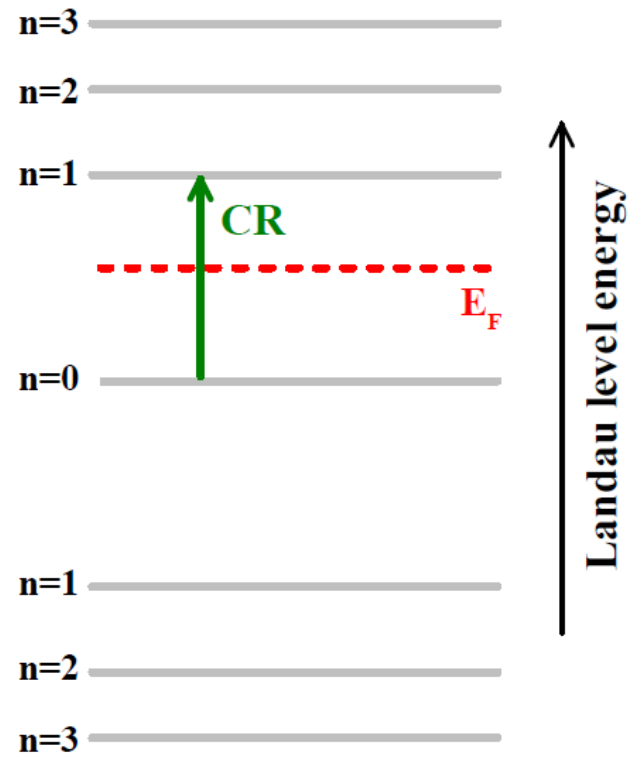
...massless yes, but not 3D Dirac

Dirac electrons – Landau level spectrum

3D Dirac electrons



2D Dirac electrons

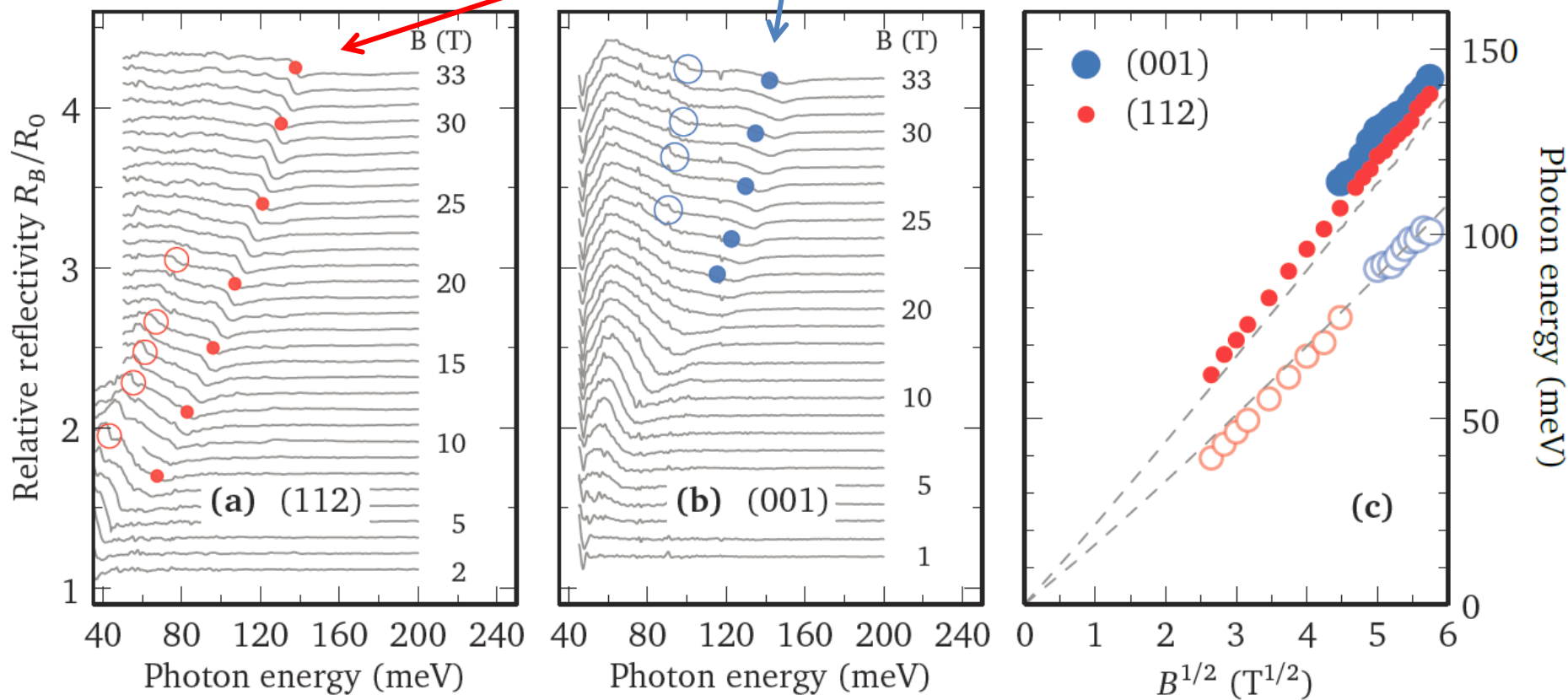


$$E_n = \pm v \sqrt{2e\hbar B n + \hbar^2 k^2}$$

$$E_n = \pm v \sqrt{2e\hbar B n}$$

Cd₃As₂ – High-field magneto-reflectivity

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Magneto-optical response linear in \sqrt{B} = typical signature of massless particles

...massless yes, but not 3D Dirac

Bodnar's model of electronic bands in Cd_3As_2

BAND STRUCTURE OF Cd_3As_2 FROM SHUBNIKOV-de HAAS
AND de HAAS - van ALPHEN EFFECTS

J. BODNAR

Department of Solid State Physics, Polish Academy of Sciences,
Zabrze, Poland

Experimental values of SdH and dHvA periods and cyclotron effective masses found by Rosenman and Doi et al. have been compared with the theoretical predictions derived in this work for a tetragonal narrow gap semiconductor. By the least square fit method the values of band parameters were obtained. It has been established that Cd_3As_2 has inverted band structure resembling HgTe under tensile stress.

Electronic bands in Cd_3As_2

Two energy scales of conical bands:

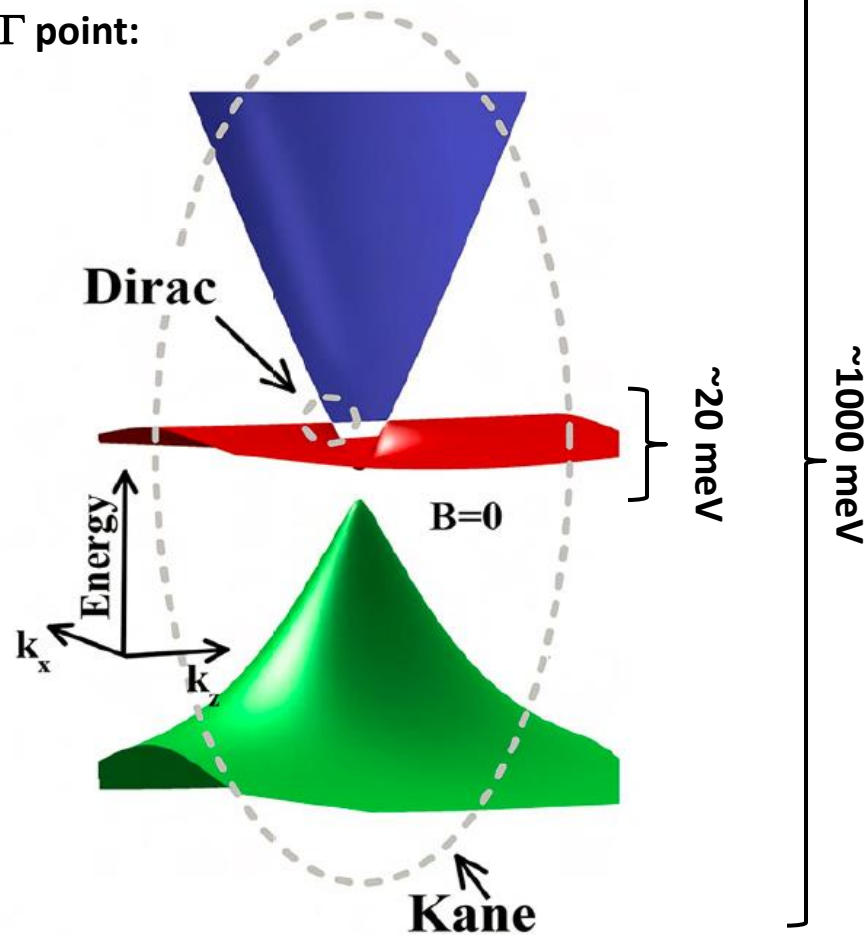
Two symmetry-protected **Dirac** cones at low energies

...due crossing of heavy and light hole band in a tetragonally distorted zinc-blende semiconductor

A single cone of massless **Kane** electrons, no symmetry protection

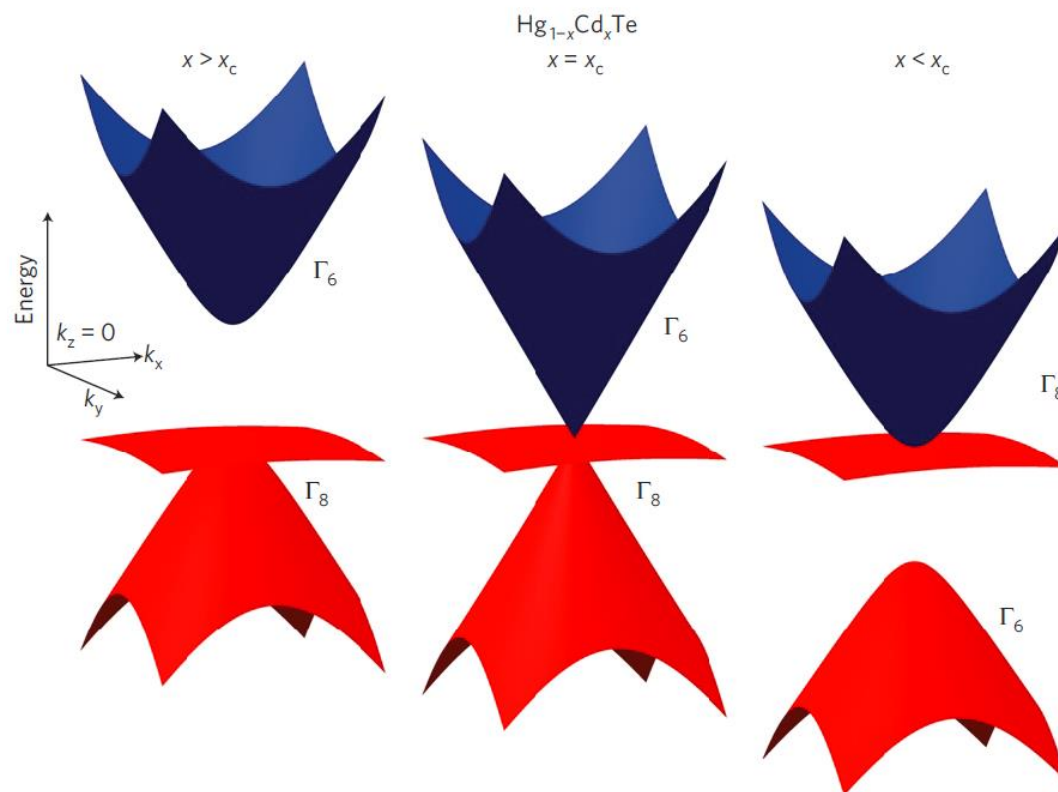
...appearing in zinc-blende semiconductors with a vanishing band gap

Γ point:



J. Bodnar, in Proc. III Conf. Narrow-Gap Semiconductors, Warsaw,(Elsevier, 1977) p. 311

Massless Kane electrons in gapless HgCdTe

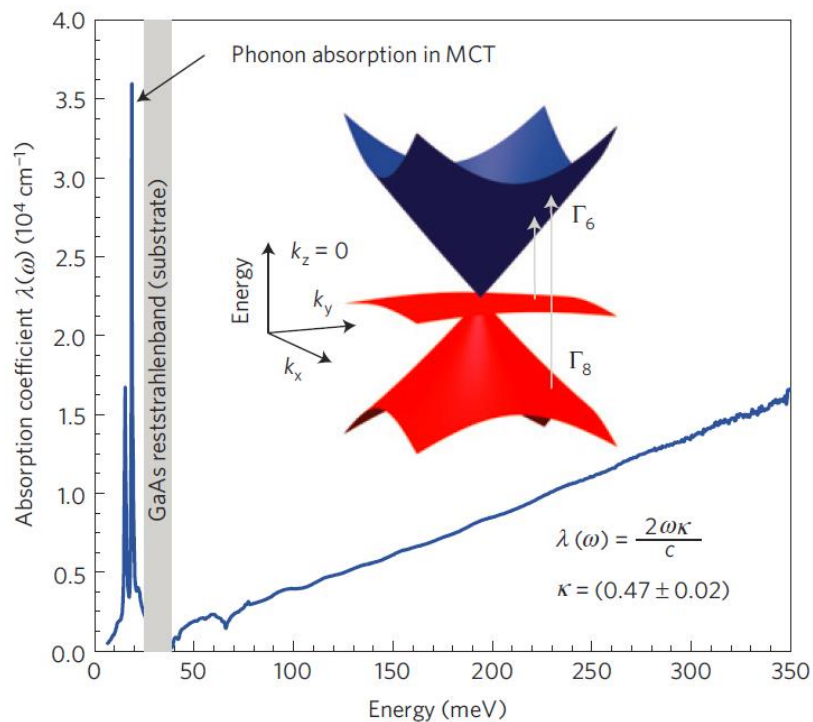


E. O. Kane, J. Phys. Chem. Solids 1, 249 (1957)
 P. Kacman and W. Zawadzki, phys. stat. sol. (b) 47, 629 (1971)
 M. Orlita et al., Nature Phys. 10, 233 (2014)

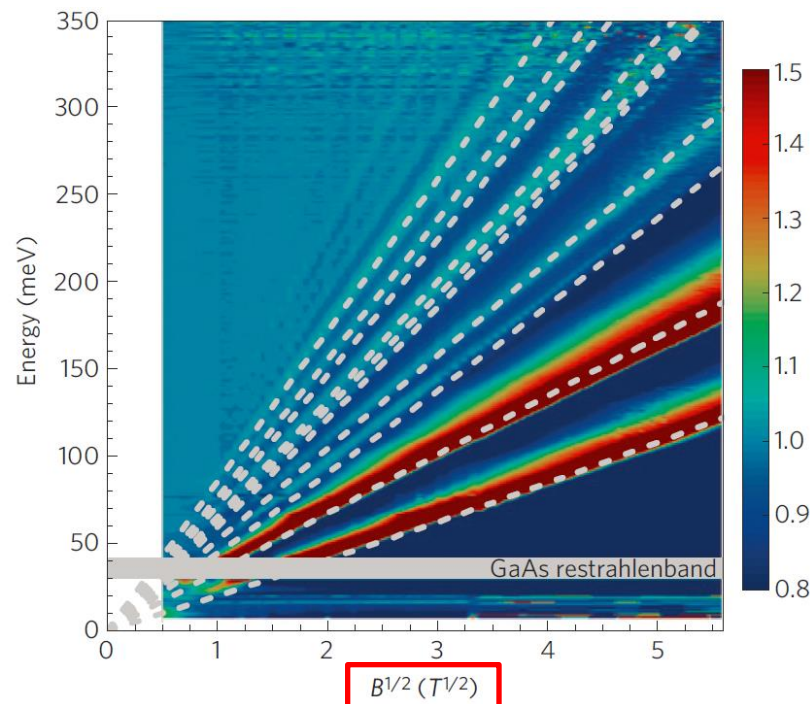
A single conical band in the center of Brillouin zone (due to accidental degeneracy of levels) hosting massless Kane electrons

Massless Kane electrons in gapless HgCdTe

Absorption coefficient linear in photon energy:



Magneto-optical response linear in \sqrt{B} :



Electronic bands in Cd_3As_2

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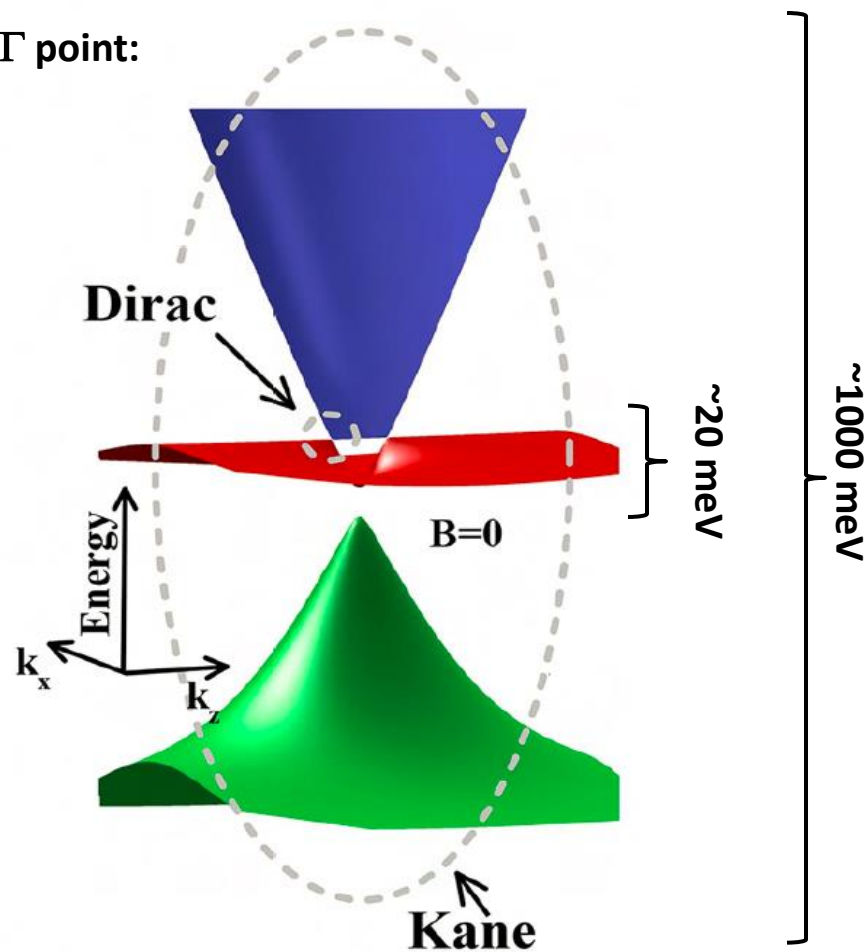
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A. Akrap et al., arXiv:1604.00038 (2016)

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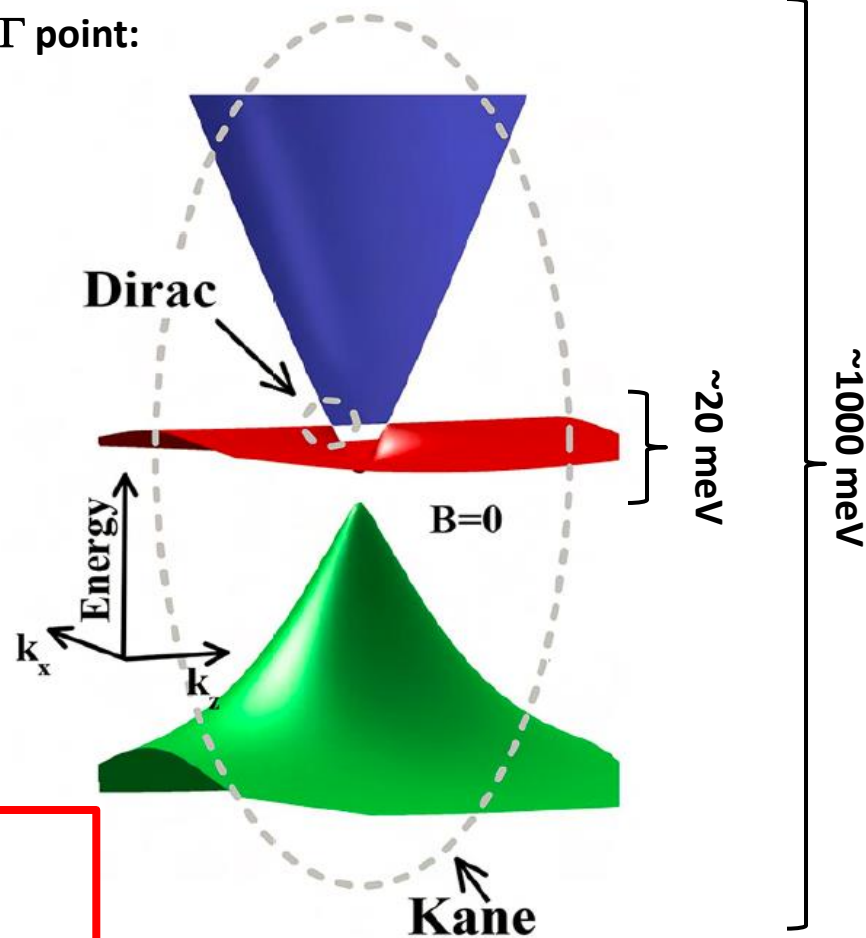
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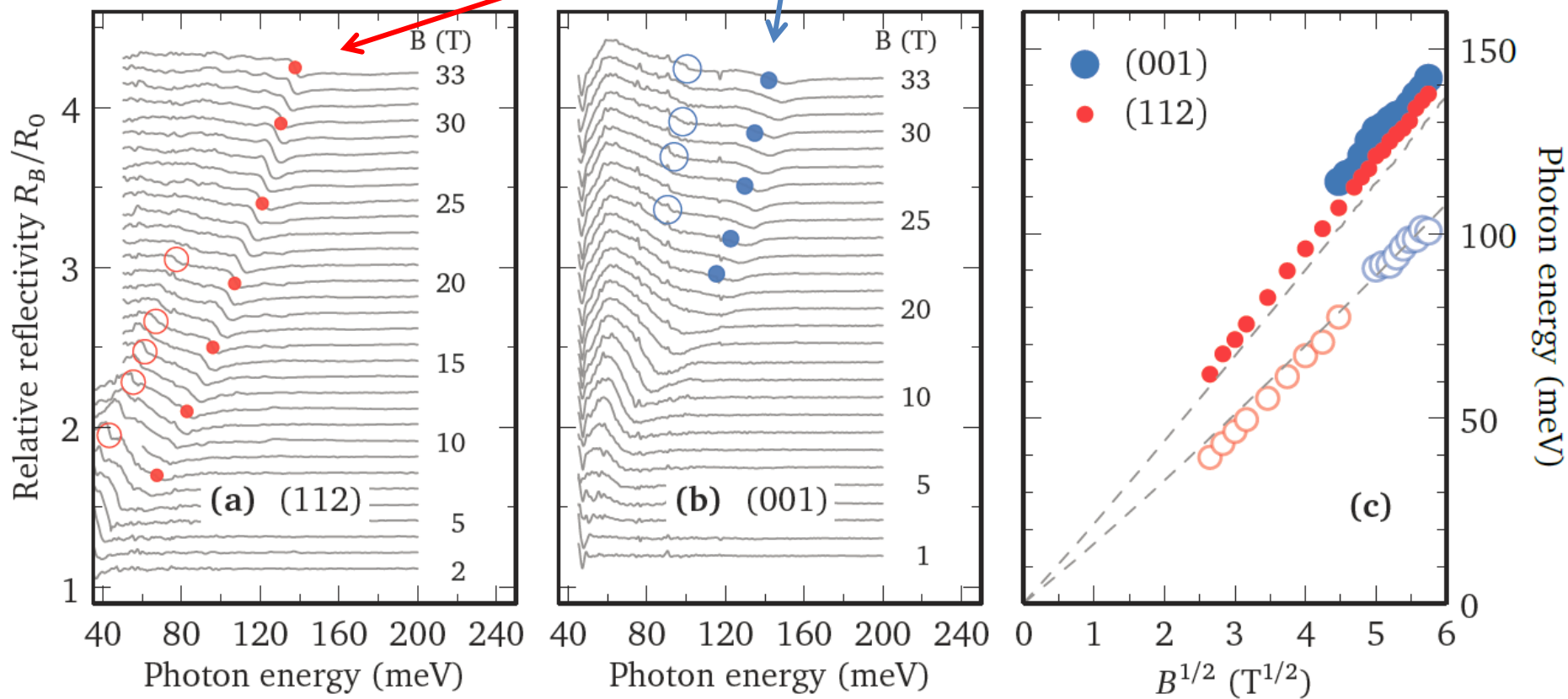
Γ point:



A. Akrap, Wednesday 1/6/16, 9:50

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Cyclotron resonance (CR) in the quantum limit

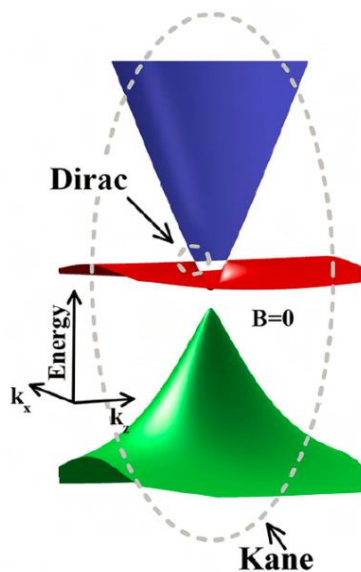
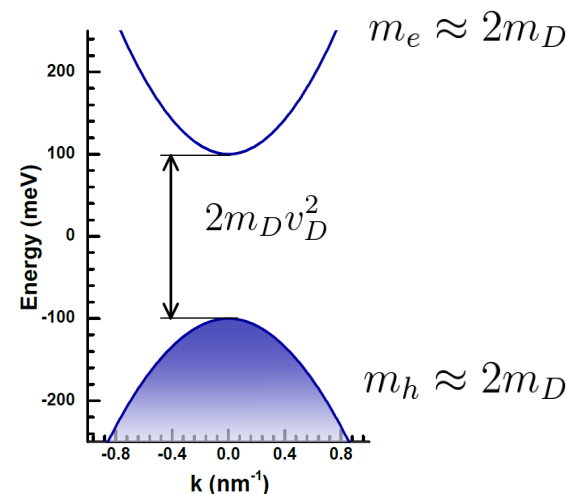


Magneto-optical signature of **massless Kane electrons**

Conclusions

Electrons and holes in bulk Bi_2Se_3 closely resemble massive Dirac particles in quantum electrodynamics

M. Orlita et al., Phys. Rev. Lett. 114, 186401 (2015)



The band structure of Cd_3As_2 hosts two kinds of 3D conical features:

3D massless Dirac and Kane electrons at a “small” and “big” energy scale, respectively

A. Akrap et al., arXiv:1604.00038 (2016)

